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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER ENVPREDRSCHFAC Technical Paper No. 5-74	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) A Three-Parameter Model for Limited Area Forecasting		5. TYPE OF REPORT & PERIOD COVERED
		6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(s) Dr. L. Bengtsson Swedish Meteorological and Hydrological Institute, Stockholm, Sweden		8. CONTRACT OR GRANT NUMBER(s)
9. PERFORMING ORGANIZATION NAME AND ADDRESS Environmental Prediction Research Facility Naval Postgraduate School Monterey, California 93940		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
11. CONTROLLING OFFICE NAME AND ADDRESS Commander, Naval Air Systems Command Department of the Navy Washington, D.C. 20361		12. REPORT DATE March 1974
		13. NUMBER OF PAGES 112
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		15. SECURITY CLASS. (of this report) UNCLASSIFIED
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; Distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Atmosphere Atmospheric Prediction Numerical Limited Area Forecasting Atmospheric Circulation		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This report describes an operational quasi-geostrophic three-parameter model. The original model was developed by Dr. L. Bengtsson and has been used operationally for several years at the Swedish Meteorological and Hydrological Institute. The improved model described in this report incorporates an Ekman function and the effect of the flow over mountains as well as sensible and latent heat sources. Humidity and precipitation		

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S/N 0102-014-6601

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20. (continued)

are also predicted by the model.

An additional feature is the optional inclusion in the vorticity equation of the terms which are usually considered negligible: (1) the advection of vorticity by the divergent wind, (2) the product of relative vorticity and divergence, (3) the vertical advection of vorticity, and (4) the twisting term.

The model can easily be adapted for different geographical areas with different grid lengths and also be integrated over different vertical layers. The lateral boundary values can be fixed or allowed to vary in time and thus it is possible to apply the so called nesting or telescoping technique. A special, relatively simple method to apply the nesting techniques is described.

AN (1) AD- 777 406
 FG (2) 040200
 FG (2) 120500
 CI (3) (U)
 CA (5) ENVIRONMENTAL PREDICTION RESEARCH FACILITY (NAVY)
 MONTEREY CALIF
 TI (6) A Three-Parameter Model for Limited Area Forecasting,
 TC (8) (U)
 AU (10) Bengtsson, L.
 RD (11) Mar 1974
 PG (12) 113p
 RS (14) ENVPRDRSHCHFAC-tech-paper-5-74
 RC (20) Unclassified report
 DE (23) *Weather forecasting, *Computer programs, FORTRAN,
 Numerical analysis, Atmospheric motion, Vortices,
 Humidity
 DC (24) (U)
 ID (25) *Numerical weather forecasting, FORTRAN 4 programming
 language, PROG 3P computer program, STEP 3P computer
 program, STEPEXT computer program, Primitive equations,
 Sensible heat, Latent heat
 IC (26) (U)
 AB (27) The report describes an operational quasi-geostrophic
 three-parameter model. The original model was developed
 by Dr. L. Bengtsson and has been used operationally for
 several years at the Swedish Meteorological and
 Hydrological Institute. The improved model described in
 this report incorporates an Ekman function and the
 effect of the flow over mountains as well as sensible
 and latent heat sources. Humidity and precipitation are
 also predicted by the model. (Modified author abstract)
 AC (28) (U)
 DL (33) 01
 CC (35) 407279

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Encl: (1) "A Three-Parameter Model For Limited Area Forecasting,"
ENVPREDRSCHFAC Technical Paper No. 5-74, March 1974
1. Enclosure (1) is forwarded for information.


G. D. HAMILTON

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VII.d: Item No. 5.

LIST IX.a: Item Nos. 1, 2, 3, 10, 17, 22 and 32.
IX.b: Item No. 2.

LIST X: Item Nos. 4, 8, 14 and 32.

LIST XII: Item Nos. 1, 2, 5, 6 and 7.

LIST XIII: Item Nos. AUS-1, 5 (including Meteorology Research Center), CAN-3, DEN-1, ENG-1, 5 and 14, FIN-1, FRA-1, GER-1 and 3, HGKG-1, INDIA-4 (including Meteor. Dept., New Delhi), ITL-2, JAP-1, NZEA-1, NOR-1 and 3, SOAF-1 and 2, SWE-1, 2, 3 and 4.

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Technical Paper No. 5-74

A THREE-PARAMETER MODEL FOR LIMITED AREA FORECASTING

by

DR. L. BENGTSSON

MARCH 1974



**ENVIRONMENTAL PREDICTION RESEARCH FACILITY
NAVAL POSTGRADUATE SCHOOL
MONTEREY, CALIFORNIA 93940**

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ACKNOWLEDGEMENTS

The author wishes to thank LT Byron Maxwell, USN for his careful computations and critical checking of the manuscript. The efforts of Mrs. Winona Carlisle who typed the manuscript are also gratefully acknowledged.

1. INTRODUCTION

Numerical weather prediction with the aid of primitive equations has now been performed operationally for several years (Reiser, 1969; Shuman, 1968). Forecasts for one and two days by primitive equations indicate some improvement when compared with forecasts with the quasi-geostrophic and filtered equations. However, it has not been shown in a clear and convincing way whether this is solely due to the unfiltered part of the equations or to purely numerical improvements such as improved vertical resolution, an alternating grid, and the introduction of new physical effects such as sensible heat, latent heat, radiation, etc. When the forecasts with the primitive models are extended further in time, the improvements achieved with these models seem to be more obvious.

If a quasi-geostrophic model is compared with a primitive model with a resolution of three to four vertical layers or less, it will be found that the computational time (and cost) for the forecast with the primitive model is roughly about ten times as large as for the quasi-geostrophic model, if explicit time-integration schemes are used. The grid lengths now in use for operational weather prediction are still about 300 to 400 km.

Since significant weather disturbances have dimensions which are only three to five times that size, it is obvious that the truncation errors in the computation of the horizontal finite differences are much too large.

That grid distances of 300 km and more have been used for such a long time is naturally due to insufficient computational capacity, but may also partly be due to accustomed routine. However, a decrease in the horizontal grid length to half the size implies an increase in the number of grid points by a factor of four for the same computation area.

Due to the criterion of Courant, Friedrichs and Lewy, the time step must be decreased by a factor of two in order to maintain computational stability. If the computations can be organized in the same way as for the larger grid, a halving of the grid length, therefore, implies an eight-fold increase in the computation time. From the operational point of view, therefore, a primitive model should be compared with a quasi-geostrophic model where the horizontal grid length has been decreased to half the size.

When the forecasting areas are small and cover only one third or less of a hemisphere, the horizontal boundaries will fall in meteorologically active areas. This disadvantage does not create any large problem for the filtered models and a moderate horizontal smoothing in the neighborhood of the boundaries is sufficient. For the primitive models the

problem is much worse, since high-amplitude gravity waves are generated at the boundaries. These waves propagate into the area and greatly affect the meteorological information. Except for some successful experiments (Bushby, 1967, 1968; Gerrity, 1969), there has not been reported any adequate technique to avoid this in the general case. For this reason, forecasts for restricted areas with the complete equations in operational use imply considerable difficulties.

It may now be argued that it is of no use to apply a quasi-geostrophic model to a fine mesh, since that means the model will predict (or try to predict) scales of motion characterized by large Rossby numbers. For instance, when the Rossby number is on the order of one, all the terms in the vorticity equation are of equal magnitude. The quasi-geostrophic models will thus, according to this analysis, give rise to intolerably large errors for small and intense vortices, especially at low latitudes. However, it has been shown (Bengtsson and Moen, 1971) that substantial improvements in forecasts from quasi-geostrophic models are obtained if the grid size is reduced from 300 to 150 km.

The reason for this is that the higher order terms in the vorticity equation to a considerable degree cancel each other. It is only in the final stage of the cyclone development, when the flow becomes very deformed, that these terms become important.

It is also possible, as will be shown in this report, to include these terms in the integration and estimate them with the aid of quasi-geostrophic divergence.

It is the experience of the author that filtered models still are very useful for short-range predictions at medium and high latitudes. It is also very probable that unsuccessful predictions of especially rapid cyclogenesis are due to inaccuracies in the initial state and to unsatisfactory ways of including topographical effects, parameterization of dissipation, and the heating mechanisms. Recent comparisons performed in Sweden between filtered and primitive equation models support this view.

2. PROGNOSTIC EQUATIONS

The vorticity equation, thermodynamical equation and the continuity equation read:

$$\frac{\partial \zeta}{\partial t} = - \mathbf{W} \cdot \nabla \eta - f \quad D; \quad (2.1a)$$

$$\frac{\partial}{\partial t} \left(\frac{\partial \phi}{\partial p} \right) = - \mathbf{W} \cdot \nabla \left(\frac{\partial \phi}{\partial p} \right) - R(\Gamma_d - \Gamma) \frac{\omega}{p} - \frac{R}{c_p p} \left(\frac{\delta Q}{dt} \right); \quad (2.1b)$$

$$\frac{\partial \omega}{\partial p} = - D; \quad (2.1c)$$

where

$$\mathbf{W} = \mathbf{k} \times \nabla \phi,$$

$$D = \nabla \cdot \mathbf{W},$$

$$\Gamma_d = \frac{1}{\rho c_p} \quad , \quad \text{and}$$

$$\Gamma = \frac{\partial T}{\partial p} \quad .$$

We will now introduce some special model assumptions. We assume that the atmosphere is bounded by a pressure surface near the surface of the earth, $p=p_0$, and an upper pressure surface, $p=p_1$, near the tropopause. We will further assume a third interjacent pressure surface, $p=p_m$, which separates the atmosphere in two layers. We will now represent the wind field in the following way:

$$W = W_m - 2W_1 \frac{p-p_m}{p_o-p_m} ; \quad \text{layer 1}$$

$$W = W_m + 2W_2 \frac{p_m-p}{p_m-p_1} ; \quad \text{layer 2}$$

$$W = (W_m + 2W_2) \frac{p}{p_1} ; \quad \text{layer 3}$$

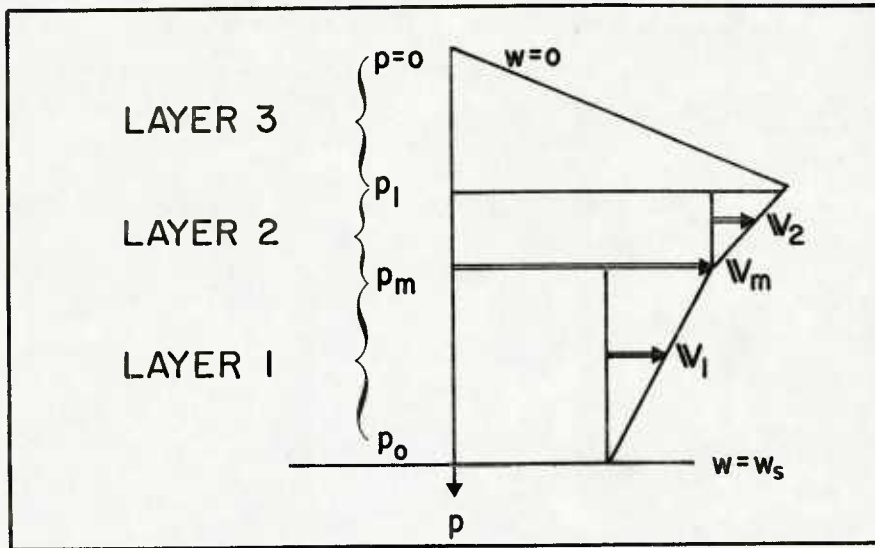


Figure 1. Vertical representation of the wind for a 3-parameter model.

According to the definition, equations for D and ζ will have the same form as those for W . Equation (2.1a) can now be integrated between $p=p_o$ and $p=0$ with the boundary conditions $\omega(p=0)=0$ and $\omega(p_o)=\omega_s$.

This gives a prognostic equation for the vertically integrated vorticity:

$$\begin{aligned} \frac{\partial}{\partial t} (\zeta_m + c_1 \zeta_2 - c_2 \zeta_1) = & -W_m \cdot \nabla (c_3 \eta_m - c_2 \zeta_1 + c_4 \zeta_2 + c_7 f) \\ & - W_1 \cdot \nabla (-c_2 \eta_m + c_5 \zeta_1) - W_2 \cdot \nabla (c_4 \eta_m + c_6 \zeta_2 + 2c_7 f) + c_8 f \omega_s \end{aligned} \quad (2.2)$$

where

$$\begin{aligned}
 c_1 &= \frac{2p_m}{2p_o - p_1} & c_4 &= \frac{6p_m - 2p_1}{6p_o - 3p_1} & c_7 &= \frac{p_1}{6p_o - 3p_1} \\
 c_2 &= \frac{2(p_o - p_m)}{2p_o - p_1} & c_5 &= \frac{8p_o - 8p_m}{6p_o - 3p_1} & c_8 &= \frac{2}{2p_o - p_1} \\
 c_3 &= \frac{6p_o - 4p_1}{6p_o - 3p_1} & c_6 &= \frac{8p_m}{6p_o - 3p_1} .
 \end{aligned}$$

The next two prognostic equations are computed from the difference between the vorticity equation at level p_m and p_o and the corresponding difference for p_1 and p_m . We will now get two vorticity equations valid for layer 1 and layer 2:

$$\frac{\partial \zeta_1}{\partial t} + (W_m - W_1) \cdot \nabla \zeta_1 + W_1 \cdot \nabla (\eta_m - \zeta_1) + f D_1 = 0; \quad (2.3)$$

$$\frac{\partial \zeta_2}{\partial t} + (W_m + W_2) \cdot \nabla \zeta_2 + W_2 \cdot \nabla (\eta_m + \zeta_2) + f D_2 = 0. \quad (2.4)$$

From these two equations we will now eliminate the divergencies D_1 and D_2 with the aid of the continuity equation (2.1c) and the thermodynamical equation (2.1b). We will first integrate the continuity equation between the two levels p_a and p_b :

$$\omega(p_a) = \omega(p_b) + \int_{p_a}^{p_b} D \, dp. \quad (2.5)$$

We will now put $p_a=p$, $p_b=p_o$ and $\omega(p_o)=\omega_s$ and will, thereby, get an expression for $\omega(p)$ in layer 1:

$$\omega(p) = \omega_s + (p_o - p) D_m + \frac{(p - p_m)^2 - (p_o - p_m)^2}{p_o - p_m} D_1. \quad (2.6)$$

With $p_a=p$, $p_b=p_m$ and $\omega(p_m)$ from (2.6) we get the following expression for ω in layer 2:

$$\omega(p) = \omega_s + (p_o - p) D_m - (p_o - p_m) D_1 + \frac{(p_m - p)^2}{p_m - p_1} D_2. \quad (2.7)$$

With $p_a=p$, $p_b=0$ and $\omega(p=0)=0$ we will have for layer 3:

$$\omega(p) = -\frac{p^2}{2p_1} (D_m + 2D_2). \quad (2.8)$$

Finally we get a relation between D_m , D_1 , D_2 and ω_s with the aid of an integration of the continuity equation from the top to the bottom of the atmosphere:

$$D_m = c_2 D_1 - c_1 D_2 - c_8 \omega_s. \quad (2.9)$$

We now introduce the stream function into the thermodynamical equation and integrate through layers 1 and 2.

$(\Gamma_d - \Gamma)$ is assumed to be constant in every layer and $\left(\frac{\delta Q}{dt}\right)_{pm} =$

$$\left(\frac{\delta Q}{dt}\right)_{p_1} = 0.$$

$$\begin{aligned}
2f \frac{\partial \psi_1}{\partial t} &= -2f \mathbf{W}_m \cdot \nabla \psi_1 + R(\Gamma_d - \Gamma)_1 \int_{p_m}^{p_o} \frac{\omega}{p} dp + \frac{R}{2c_p} \ln\left(\frac{p_o}{p_m}\right) \left(\frac{\delta Q}{dt}\right)_{p_o}; \\
2f \frac{\partial \psi_2}{\partial t} &= -2f \mathbf{W}_m \cdot \nabla \psi_2 + R(\Gamma_d - \Gamma)_2 \int_{p_1}^{p_m} \frac{\omega}{p} dp.
\end{aligned} \tag{2.10}$$

The integrals $\int \frac{\omega}{p} dp$ can be computed with the aid of equations (2.6), 2.7) and (2.9), and the system (2.10) can be written

$$\begin{aligned}
m_1 D_1 + m_2 D_2 &= H_1; \\
n_1 D_1 + n_2 D_2 &= H_2;
\end{aligned} \tag{2.11}$$

where

$$\begin{aligned}
m_1 &= \frac{R}{2} (\Gamma_d - \Gamma)_1 \left[\frac{p_1 (p_o - p_m)^2 + p_m^2 (2p_o - p_1)}{(p_o - p_m) (2p_o - p_1)} \cdot \ln\left(\frac{p_o}{p_m}\right) - \frac{2(p_o - p_m)^2}{2p_o - p_1} + \frac{1}{2} (p_o + p_m) - 2p_m \right]; \\
m_2 &= \frac{R}{2} (\Gamma_d - \Gamma)_1 \left[-\frac{2p_o p_m}{2p_o - p_1} \cdot \ln\left(\frac{p_o}{p_m}\right) + \frac{2p_m (p_o - p_m)}{2p_o - p_1} \right]; \\
n_1 &= \frac{R}{2} (\Gamma_d - \Gamma)_2 \left[\frac{p_1 (p_o - p_m)}{2p_o - p_1} \ln\left(\frac{p_m}{p_1}\right) - \frac{2(p_o - p_m) (p_m - p_1)}{2p_o - p_1} \right];
\end{aligned}$$

$$n_2 = \frac{R}{2} (\Gamma_d - \Gamma)_2 \left[\frac{p_1 p_m (2p_o - p_m)}{(2p_o - p_1)(p_m - p_1)} \ln \left(\frac{p_m}{p_1} \right) + \frac{2p_m (p_m - p_1)}{2p_o - p_1} + \frac{1}{2} (p_m + p_1) - 2p_m \right];$$

$$H_1 = f \frac{\partial \psi_1}{\partial t} + f \mathbf{V}_m \cdot \nabla \psi_1 - \frac{R}{4c_p} \ln \left(\frac{p_o}{p_m} \right) \cdot \left(\frac{\delta Q}{\delta t} \right)_{p_o} - \frac{R}{2} (\Gamma_d - \Gamma)_1 \left[- \frac{p_1}{2p_o - p_1} \ln \left(\frac{p_o}{p_m} \right) + \frac{2(p_o - p_m)}{2p_o - p_1} \right] \omega_s ;$$

$$H_2 = f \frac{\partial \psi_2}{\partial t} + f \mathbf{V}_m \cdot \nabla \psi_2 - \frac{R}{2} (\Gamma_d - \Gamma)_2 \left[- \frac{p_1}{2p_o - p_1} \ln \left(\frac{p_m}{p_1} \right) + \frac{2(p_m - p_1)}{2p_o - p_1} \right] \omega_s .$$

From the system (2.11) we can easily express the divergencies in an explicit way:

$$D_1 = -a_1 H_1 + a_2 H_2; \tag{2.12}$$

$$D_2 = b_1 H_1 - b_2 H_2;$$

where

$$a_1 = \frac{-n_2}{m_1 n_2 - n_1 m_2} ; \quad a_2 = \frac{-m_2}{m_1 n_2 - n_1 m_2} ;$$

$$b_1 = \frac{-n_1}{m_1 n_2 - n_1 m_2} ; \quad b_2 = \frac{-m_1}{m_1 n_2 - n_1 m_2} .$$

Introducing these expressions for D_1 and D_2 into the prognostic equations (2.3) and (2.4) and expressing the wind and vorticity in terms of the stream function gives:

$$\begin{aligned} \nabla^2 \frac{\partial \psi_1}{\partial t} - a_1 f^2 \frac{\partial \psi_1}{\partial t} + a_2 f^2 \frac{\partial \psi_2}{\partial t} = & -J(\psi_m; \zeta_1) - J(\psi_1; n_m - 2\zeta_1) \\ & + a_1 f^2 J(\psi_m; \psi_1) - a_2 f^2 J(\psi_m; \psi_2) - a_3 f \omega_S - a_1 f H \end{aligned} \quad (2.13)$$

$$\begin{aligned} \nabla^2 \frac{\partial \psi_2}{\partial t} - b_2 f^2 \frac{\partial \psi_2}{\partial t} + b_1 f^2 \frac{\partial \psi_1}{\partial t} = & -J(\psi_m; \zeta_2) - J(\psi_2; n_m + 2\zeta_2) \\ & - b_1 f^2 J(\psi_m; \psi_1) + b_2 f^2 J(\psi_m; \psi_2) + b_3 f \omega_S + b_1 f H \end{aligned} \quad (2.14)$$

Here we have

$$a_3 = a_1 s_1 - a_2 s_2;$$

$$b_3 = b_1 s_1 - b_2 s_2;$$

where

$$\begin{aligned} s_1 &= \frac{R}{2} (\Gamma_d - \Gamma)_1 \left[-\frac{p_1}{2p_o - p_1} \ln\left(\frac{p_o}{p_m}\right) + \frac{2(p_o - p_m)}{2p_o - p_1} \right]; \\ s_2 &= \frac{R}{2} (\Gamma_d - \Gamma)_2 \left[-\frac{p_1}{2p_o - p_1} \ln\left(\frac{p_m}{p_1}\right) + \frac{2(p_m - p_1)}{2p_o - p_1} \right]; \\ H &= \frac{R}{4c_p} \ln\left(\frac{p_o}{p_m}\right) \left(\frac{\delta Q}{\delta t}\right)_{p_o} = h_1 \left(\frac{\delta Q}{\delta t}\right)_{p_o} \text{ with } h_1 = \frac{R}{4c_p} \ln\left(\frac{p_o}{p_m}\right). \end{aligned}$$

We will also introduce the stream functions into equation (2.2) and define ψ_M by $\psi_M = \psi_m - c_2 \psi_1 + c_1 \psi_2$.

This results in the equation

$$\nabla^2 \frac{\partial \psi_M}{\partial t} = - J(\psi_m; c_3 n_m - c_2 \zeta_1 + c_4 \zeta_2 + c_7 f) - J(\psi_1; -c_2 n_m + c_5 \zeta_1) \\ - J(\psi_2; c_4 n_m + c_6 \zeta_2 + 2c_7 f) + c_8 f \omega_s. \quad (2.15)$$

The equations (2.13) through (2.15) now constitute our system of prognostic equations. The only thing which we now have to do is to find an expression for the non-adiabatic heat H , and the lower boundary condition, ω_s .

3. BOUNDARY CONDITIONS

3.1 LOWER BOUNDARY CONDITIONS

We now assume that ω_s can be separated into two parts; one part which depends upon the dissipation in the boundary layer ω_s' and one part which depends upon topography ω_s'' . We thereby assume:

$$\omega_s = \omega_s' + \omega_s'' . \quad (3.1.1)$$

From the Ekman theory about the variation of the wind in the friction layer, we could easily derive the following expression:

$$\omega_s' = - g \rho_o \sqrt{\frac{K}{2f}} \cdot F \cdot \zeta_o \quad \text{where } F = 1 + c \cdot \sin \theta - c \cdot \cos \theta,$$

the wind in the surface layer has a magnitude $c|V_o|$, the angle between the geostrophic wind and the surface wind is given by θ , K is the turbulent coefficient of the viscosity, and ρ_o is the density of the air at the surface. Inserting the following numerical values:

$$K = 10 \text{ m/s}; \quad g = 9.81 \text{ m/s}^2; \quad f = 10^{-4} \text{ s}^{-1}; \quad T_o = 280^\circ \text{K};$$

$$R = 287; \quad p_o = 100 \text{ cb}; \quad \text{gives}$$

$$\omega_s' = - 2.729597 F \cdot \zeta_o = -2.729597 \cdot F(\zeta_m - 2\zeta_1). \quad (3.1.2)$$

$F(c, \theta)$ can be given a constant value in the model or we can also assume different values over land and sea.

The following values will be used:

Over land $\theta = 10^\circ$ and $c = 0.78$ gives $F = 0.36635$,

that is $\omega'_s = -1. \zeta_o$.

Over sea $\theta = 5^\circ$; $c = 0.85$ gives $F = 0.22734$,

that is, $\omega'_s = 0.62055 \zeta_o$.

If k_2 is assumed to be that part of the air which is forced over the mountains, we get

$$\omega''_s = k_2 \nabla_o \cdot \nabla p_s = k_2 (\nabla_m - 2\nabla_1) \cdot \nabla p_s \quad (3.1.3)$$

$k_2 = 1$ will be used in the model.

The equations (3.1.1 - 3.1.3) now give the lower boundary condition for ω_s . For further information see Bengtsson (1969). This way of treating the topographical effect as given by equation 3.1.3 is quite unrealistic and seems to underestimate the effect of the topography. (This will be especially true for steep and/or small scale mountains.) A new way to treat mountains as impenetrable vertical barriers has recently been published (Egger, 1972). A similar way to include mountains in vertically integrated balanced models will be described by the author in a coming investigation.

3.2 LATERAL BOUNDARY CONDITIONS

The model can use two different kinds of lateral boundary conditions namely:

a. Constant Inflow

$$\frac{\partial \psi_1}{\partial t} = \frac{\partial \psi_2}{\partial t} = \frac{\partial \psi_m}{\partial t} = 0$$

$$\zeta_1(t) = \zeta_1(t=0) \quad (= \text{constant in time})$$

$$\zeta_2(t) = \zeta_2(t=0) \quad (= \text{constant in time}) \quad (3.2.1)$$

$$\zeta_m(t) = \zeta_m(t=0) \quad (= \text{constant in time})$$

b. Variable Inflow

The values for $\frac{\partial \psi_1}{\partial t}$, $\frac{\partial \psi_2}{\partial t}$, $\frac{\partial \psi_m}{\partial t}$, ζ_1 , ζ_2 , and ζ_m along the boundary are generated through an integration for a larger area which includes the actual area. Interpolation in time and space is necessary to synchronize the boundary values. A technical description of this is found in section 10.

4. PARAMETERIZATION OF PHYSICAL PROCESSES

4.1 SENSIBLE HEAT

In the computation of equation(2.10) we assume that $\frac{\delta Q}{dt}$ decreased linearly to 0 from p_o to p_m . The sensible heat which is transported to the atmosphere from the underlying surface, is introduced in the following way:

$$\left(\frac{\delta Q}{dt}\right)_{p_o} = (A_1 |V_o| + A_2) (T_s - T_o) \text{ over sea, if } T_s > T_o; \quad (4.1.1)$$

$$\left(\frac{\delta Q}{dt}\right)_{p_o} = 0 \text{ over land and over ocean areas if } T_s \leq T_o;$$

where A_1 and A_2 are empirical constants $A_2 = 10A_1$
 $= 5 \cdot 10^{-3} \text{ m(sec)}^{-2} (\text{deg})^{-1}$. T_s is the sea surface temperature and T_o is the air temperature near the sea surface.

An approximative temperature at the level $p_x = \frac{1}{2} (p_o + p_m)$ is obtained through an integration of the hydrostatic equation:

$$T_x = \frac{2f_o}{R \ln \left(\frac{p_o}{p_m}\right)} \psi_1. \quad (4.1.2)$$

We now assume that Γ is constant in the layer and is 75% of Γ_d . We therefore get

$$T_o = T_x + 0.75 \Gamma_d \frac{1}{2} (p_o - p_m) = \frac{2f_o}{R \ln \left(\frac{p_o}{p_m}\right)} \left[1 + \frac{0.75 R (p_o - p_m)}{c_p (p_o + p_m)} \right] \psi_1. \quad (4.1.3)$$

In equations (2.13) and (2.14) the sensible heating (H) was defined to be:

$$H = \frac{R}{4C_p} \ln \left(\frac{p_o}{p_m} \right) \left(\frac{\delta Q}{dt} \right)_{p_o} .$$

Substitution of $\left(\frac{\delta Q}{dt} \right)_{p_o}$ from equation (4.1.1) and using the expression for T_o from equation (4.1.3) gives:

$$H = h_1 \cdot 10^{-2} (0.5 \cdot 10^{-1} |v_o| - 0.5) (T_s - h_2 \cdot \psi_1) \quad (4.1.4)$$

where

$$h_1 = \frac{R}{4c_p} \ln \left(\frac{p_o}{p_m} \right) ;$$

$$h_2 = \frac{2f_o}{R \ln \left(\frac{p_o}{p_m} \right)} \left[1 + \frac{0.75 R(p_o - p_m)}{c_p (p_o + p_m)} \right] .$$

4.2 PROGNOSTIC EQUATION FOR HUMIDITY

The specific humidity can be predicted by the following equation:

$$\frac{\partial q}{\partial t} = - \nabla \cdot \nabla q - \omega \frac{\partial q}{\partial p} + \epsilon - r + A \nabla^2 q . \quad (4.2.1)$$

Here ϵ denotes evaporation, r condensation, and A is the coefficient of dissipation. We will disregard ϵ and r will be introduced in a different way. Using the continuity equation we get the following prognostic equation:

$$\frac{\partial q}{\partial t} = - \nabla \cdot (q \nabla) - \frac{\partial}{\partial p} (q \omega) + A \nabla^2 q . \quad (4.2.2)$$

Since our model has a very low vertical resolution we have to parameterize the vertical distribution of humidity. We will, therefore, use precipitable water as the prognostic variable. The precipitable water is defined as

$$p_w = \int_0^{\infty} \rho_w dz = \frac{1}{g} \int_{p_T}^{p_O} q dp \quad (4.2.3)$$

where

ρ_w = the density of water vapor and p_T is the pressure at the level over which we can disregard the humidity. (Here we have $p_T = 30$ cb and $p_O = 100$ cb.)

A new quantity

$$w = \frac{1}{p_O - p_T} p_w \quad (4.2.4)$$

which may be called normalized precipitable water is introduced. We assume that $q(x,y,p,t) = gE(p)w(x,y,t)$ where $E(p)$ describes the vertical variation of q computed from the standard atmosphere under the assumption that the relative humidity is 50%. The expression (4.2.4) is now introduced into (4.2.2) and we then integrate with respect to p from p_O to p_1 to give

$$\frac{\partial w}{\partial t} = - \nabla \cdot (\tilde{V} w) - w d' \omega_s + A \nabla^2 w \quad (4.2.5)$$

where

$$\tilde{V} = \frac{1}{p_O - p_T} \int_{p_T}^{p_O} V(p) E(p) dp \text{ and } d' = \frac{E(p_O)}{p_O - p_T} .$$

With $\nabla = kx\nabla\psi + \nabla\chi$ equation (4.2.5) can be written

$$\frac{\partial w}{\partial t} = -J(\tilde{\psi}, w) - \nabla\tilde{\chi} \cdot \nabla w - w\nabla^2\tilde{\chi} - wd'\omega_s + A\nabla^2 w \quad (4.2.6)$$

where

$$\tilde{\psi} = e_m\psi_m + e_1\psi_1 + e_2\psi_2 \quad \text{and} \quad \nabla^2\tilde{\chi} = \tilde{D} = e_mD_m + e_1D_1 + e_2D_2.$$

4.3 LATENT HEAT

There are two conditions for condensation:

- (a) $\bar{\omega}_1 < \omega_{tol}$ (where ω_{tol} is a given tolerance)
- (b) the relative humidity should exceed and be equal to 80%.

If these two conditions are valid, the latent heat is computed by the aid of expression (4.3.1)

$$H_{lat} = \frac{R}{4c_p} \ln \left(\frac{p_o}{p_m} \right) \left(\frac{\delta Q}{dt} \right)_{lat} = h_1 \left(\frac{\delta Q}{dt} \right)_{lat}$$

where

(4.3.1)

$$\begin{aligned} \left(\frac{\delta Q}{dt} \right)_{lat} = & -L \cdot \bar{\omega}_1 F \quad \text{and} \quad F = \frac{1}{1 + \frac{L}{c_p} \left(\frac{\partial q^x}{\partial t} \right)_p} \left[\left(\frac{\partial q^x}{\partial p} \right)_T \right. \\ & \left. + \frac{R}{c_p} \frac{T}{p} \left(\frac{\partial q^x}{\partial T} \right)_p \right] \end{aligned}$$

q^x is the maximum specific humidity. If the conditions (a) and (b) are not valid, we put $H_{lat} = 0$. Also see paragraph 8.4.

5. COMPUTATION OF THE VERTICAL MOTION

The integrated vertical motions in the two layers 1 and 2 are computed in the model. With the aid of the equation (2.6) and (2.9) we obtain:

$$\begin{aligned}\bar{\omega}_1 &= \frac{1}{p_o - p_m} \int_{p_m}^{p_o} \omega dp = \frac{p_o + p_m - p_1}{2p_o - p_1} \omega_s + \frac{\frac{1}{3}(2p_1 - 3p_m - p_o)(p_o - p_m)}{2p_o - p_1} D_1 \\ &\quad - \frac{p_m(p_o - p_m)}{2p_o - p_1} D_2;\end{aligned}\quad (5.1)$$

$$\bar{\omega}_1 = t_1 \omega_s + t_2 D_1 - t_3 D_2.$$

In the same way we obtain from the equation (2.7) and (2.9):

$$\begin{aligned}\bar{\omega}_2 &= \frac{1}{p_m - p_1} \int_{p_1}^{p_m} \omega dp = \frac{p_m}{2p_o - p_1} \omega_s - \frac{p_m(p_o - p_m)}{2p_o - p_1} D_1 + \\ &\quad + \frac{\frac{1}{3}(2p_o - p_1)(p_m - p_1) - p_m(2p_o - p_m - p_1)}{2p_o - p_1} D_2;\end{aligned}\quad (5.2)$$

$$\bar{\omega}_2 = t_4 \omega_s - t_3 D_1 + t_5 D_2.$$

Here we have

$$t_1 = \frac{p_o + p_m - p_1}{2p_o - p_1}$$

$$t_2 = \frac{1}{3} \frac{(2p_1 - 3p_m - p_o)(p_o - p_m)}{2p_o - p_1}$$

$$t_3 = \frac{p_m(p_o - p_m)}{2p_o - p_1}$$

$$t_4 = \frac{p_m}{2p_o - p_1}$$

$$t_5 = \frac{\frac{1}{3}(2p_o - p_1)(p_m - p_1) - p_m(2p_o - p_m - p_1)}{(2p_o - p_1)}$$

The physical parameters are computed by the subroutine COEFF3P.

(See appendix B to this report)

6. NUMERICAL VALUES OF THE CONSTANTS

For the levels $p_0 = 100$ cb, $p_m = 50$ cb, $p_1 = 30$ cb and the stabilities $(\Gamma_d - \Gamma)_1 = 0.422222$, $(\Gamma_d - \Gamma)_2 = 0.511111$ we get the following numerical values for the constants:

$$\begin{aligned}
 c_1 &= \frac{10}{17} = 0.588235 & m_1 &= -826.330 & s_1 &= 28.230856 \\
 c_2 &= \frac{10}{17} = 0.588235 & m_2 &= -688.40 & s_2 &= 10.645832 \\
 c_3 &= \frac{48}{51} = 0.941176 & n_1 &= -532.92 & a_3 &= 0.044374 \\
 c_4 &= \frac{24}{51} = 0.470588 & n_2 &= -10.58.36 & b_3 &= 0.012258 \\
 c_5 &= \frac{40}{51} = 0.784314 & a_1 &= 2.0828 \cdot 10^{-3} & h_1 &= 0.495351 \cdot 10^{-1} \\
 c_6 &= \frac{40}{51} = 0.784314 & a_2 &= 1.3546 \cdot 10^{-3} & h_2 &= 0.110953 \cdot 10^{-5} \\
 & & & & h_3 &= 1.03552 \cdot 10^{-6} \\
 & & & & h_4 &= 0 \\
 & & & & h_5 &= 0.492929 \cdot 10^{-1} \\
 & & & & h_6 &= 1.03552 \cdot 10^{-6} \\
 c_7 &= \frac{3}{51} = 0.588235 & b_1 &= 1.0475 \cdot 10^{-3} & t_1 &= 0.705882 \\
 c_8 &= \frac{2}{170} = 0.0117647 & b_2 &= 1.6261 \cdot 10^{-3} & t_2 &= -18.627500 \\
 & & & & t_3 &= 14.705900 \\
 & & & & t_4 &= 0.294118 \\
 & & & & t_5 &= -28.627500
 \end{aligned}$$

7. INTEGRATION OF THE COMPLETE VORTICITY EQUATION

7.1 GENERAL ASPECTS

Very little knowledge exists about the effect of the small order terms in the vorticity equation: the advection of vorticity by the divergent wind, $\mathbf{W}_\chi \cdot \nabla \zeta$; the product of relative vorticity and divergence, $\zeta \nabla \cdot \mathbf{W}$; the vertical advection of vorticity, $\omega \frac{\partial \zeta}{\partial p}$; and the twisting term $\mathbf{k} \cdot \left(\frac{\partial \mathbf{W}}{\partial p} \times \nabla \omega \right)$.

If these terms are used it is necessary, in order to conserve the total energy for an adiabatic model, to use the complete balance equation. If this is not the case, the model will not conserve total energy and after a certain time the development starts to deteriorate. This judgment has been mainly qualitative and we do not know the size of the error due to this inconsistency. It may be that this error is relatively small in comparison to other errors, as for instance, uncertainties of the initial state, and uncertainties in the description of the dissipation and the heating mechanisms. Therefore, it is necessary to perform a more detailed study of the problem and base our decision on a quantitative investigation.

It is by no means evident that we should use the same kind of assumptions in the formulation of models for short-range predictions (24 hours) as for models for medium- and long-range prediction. An example will illustrate this.

If one is interested in long-time integrations of the barotropic vorticity equation, it is necessary to use a finite difference expression which conserves kinetic energy, vorticity, and mean-squared vorticity. A finite difference expression which conserves these identities has been derived by Arakawa. However, if one is interested in short-range predictions, the so-called "Arakawa Jacobian" is not recommended since the phase-speed error is larger than the conventional finite difference analog to the Jacobian operator which only conserves vorticity. Experiments have shown that for forecasts up to four or five days it is not necessary to use an energy consistent Jacobian operator since the error in the kinetic energy is much smaller than errors due to other effects.

One of the problems which we have with the simplified vorticity equation is the over-prediction of anticyclogenesis. This seems due mainly to the lack of the term $\zeta \nabla \cdot \mathbf{W}$ in the vorticity equation:

$$\frac{\partial \zeta}{\partial t} = \dots - (f + \zeta) \nabla \cdot \mathbf{W}. \quad (7.1.1)$$

In areas of convergence, relative vorticity is mostly positive, or will be after a short time. This means that the relative vorticity will increase faster in such areas if the complete expression (7.1.1) is included. On the other hand, in areas of divergence, the relative vorticity is mostly negative, or

will be after some hours. If the complete expression is used, the relative vorticity will decrease more slowly than if we use the simplified expression. It is easily seen that this term will create an asymmetry in the vorticity pattern which is also observed in reality.

Also the vertical advection of vorticity seems to play an important role, especially in cyclone development. During the development of the cyclone the activity is mainly concentrated in two different areas.

One area is in front of the warm front or, later in the development, the occluded part of the front. The other area is found below the upper-air low or trough, where a special center of activity is created in the later stages of the cyclone development. During the development of the cyclone, an area of sinking motion is concentrated under the upper-air low.

$$\frac{\partial \zeta}{\partial t} = -\omega \frac{\partial \zeta}{\partial p} \quad (7.1.2)$$

It is easily seen from the vorticity equation (7.1.2) that this effect will give an increase in the relative vorticity in areas of sinking motion and where the vorticity increases with height.

Upward motion over a surface low yields, in the same way, an increase in the vorticity for the levels above. Therefore the vertical advection of vorticity will increase the speed of occlusion. Figure 2 shows two different 12-hour

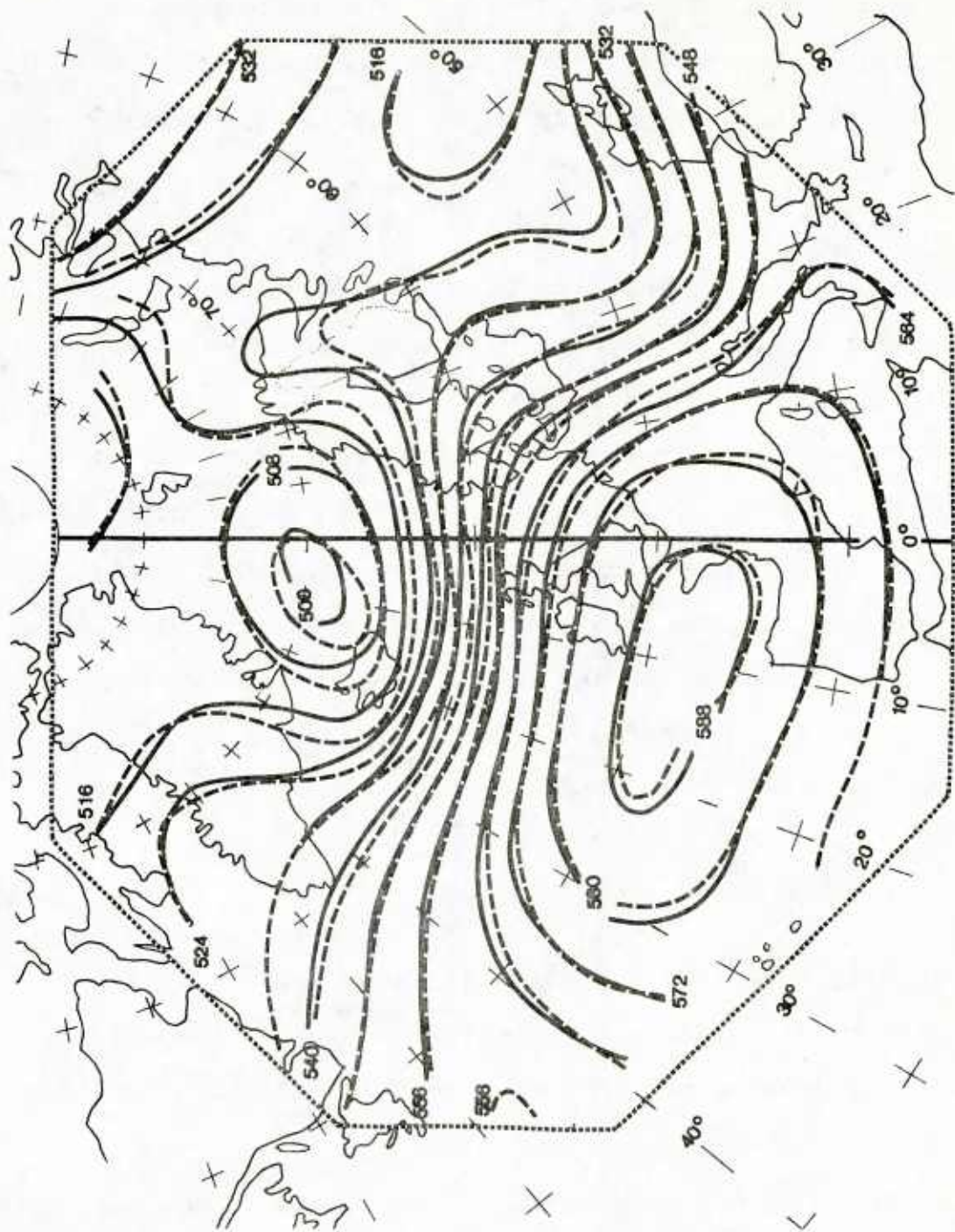


Figure 2. Two example 12-hour predictions with a 5-layer quasi-geostrophic model.

predictions with a 5-layer quasi-geostrophic model. Solid lines indicate a prediction performed with a quasi-geostrophic energy consistent model. Dashed lines indicate a prediction where the terms $-(\zeta \nabla \cdot \mathbf{V} + \omega \frac{\partial \zeta}{\partial p})$ have been included.

7.2 DERIVATION OF THE FORECASTING EQUATIONS

The complete vorticity equation reads:

$$\frac{\partial \zeta}{\partial t} = - \underbrace{\mathbf{V}_\psi \cdot \nabla \eta}_1 - \underbrace{\mathbf{V}_\chi \cdot \nabla \eta}_2 - f \nabla \cdot \mathbf{V}_\chi - \underbrace{\zeta \nabla \cdot \mathbf{V}_\chi}_3 - \underbrace{\omega \frac{\partial \zeta}{\partial p}}_4 + \underbrace{|\mathbf{k} \cdot (\frac{\partial \mathbf{V}_\psi}{\partial p} \times \nabla \omega))}_4. \quad (7.2.1)$$

Here \mathbf{V}_ψ is the non-divergent wind and \mathbf{V}_χ the divergent wind. The terms 1, 2, 3 and 4 are denoted non-geostrophic terms. Integrating through layer 1 by the representation of \mathbf{V} , ζ and D given in section 2 yields

$$\begin{aligned} (p_0 - p_m) \left[\frac{\partial \zeta_m}{\partial t} - \frac{\partial \zeta_1}{\partial t} \right] = (p_0 - p_m) \left[- \mathbf{V}_m \cdot \nabla (\zeta_m - \zeta_1 + f) + \mathbf{V}_1 \cdot \nabla (\zeta_m + f) \right. \\ \left. - \frac{4}{3} \mathbf{V}_1 \cdot \nabla \zeta_1 - f(D_m - D_1) - \zeta_m(D_m - D_1) + \zeta_1(D_m - \frac{4}{3}D_1) \right] \\ + 2\zeta_1 \bar{\omega}_1 - 2 |\mathbf{k} \cdot (\mathbf{V}_1 \times \nabla \bar{\omega}_1)| \end{aligned} \quad (7.2.2)$$

Integration through layer 2 and layer 3 yields in a similar way

$$\begin{aligned} (p_m - p_1) \left[\frac{\partial \zeta_m}{\partial t} + \frac{\partial \zeta_2}{\partial t} \right] = (p_m - p_1) \left[- \mathbf{V}_m \cdot \nabla (\zeta_m + \zeta_2 + f) - \mathbf{V}_2 \cdot \nabla (\zeta_m + f) \right. \\ \left. - \frac{4}{3} \mathbf{V}_2 \cdot \nabla \zeta_2 - f(D_m + D_2) - \zeta_m(D_m + D_2) - \zeta_2(D_m + \frac{4}{3}D_2) \right] \\ + 2\zeta_2 \bar{\omega}_2 - 2 |\mathbf{k} \cdot (\mathbf{V}_2 \times \nabla \bar{\omega}_2)| \end{aligned} \quad (7.2.3)$$

$$\begin{aligned} \frac{1}{2}p_1 \left[\frac{\partial \zeta_m}{\partial t} + 2 \frac{\partial \zeta_2}{\partial t} \right] = \frac{1}{2}p_1 \left[- (V_m + 2V_2) \cdot \nabla f - \frac{2}{3} (V_m + 2V_2) \cdot \nabla (\zeta_m + 2\zeta_2) \right. \\ \left. - f(D_m + 2D_2) - \frac{2}{3}(\zeta_m + 2\zeta_2)(D_m + 2D_2) \right] - (\zeta_m + 2\zeta_2) \bar{\omega}_3 \\ + |k \cdot \{ V_m + 2V_2 \} \times \nabla \bar{\omega}_3 \} \end{aligned} \quad (7.2.4)$$

where $\bar{\omega}_1$ and $\bar{\omega}_2$ are given in (5.1) and (5.2) and $\bar{\omega}_3$ is given by:

$$\bar{\omega}_3 = \frac{1}{p_1} \int_0^{p_1} \omega dp. \quad (7.2.5)$$

If we now add (7.2.2), (7.2.3) and (7.2.4) and divide by $\frac{2p_0 - p_1}{2}$, we get a prognostic equation for the vertically integrated mean vorticity. From now on we will only keep the non-geostrophic terms on the right hand side of the equation.

$$\begin{aligned} \frac{\partial}{\partial t}(\zeta_m + c_1 \zeta_2 - c_2 \zeta_1) = \{ - (V_\chi)_m \cdot \nabla (c_3 \eta_m - c_2 \zeta_1 + c_4 \zeta_2 + c_7 f) \\ - (V_\chi)_1 \cdot \nabla (-c_2 \zeta_m + c_5 \zeta_1) - (V_\chi)_2 \cdot \nabla (c_4 \eta_m + c_6 \zeta_2 + 2c_7 f) \}_1 \\ + \{ c_8 \zeta_m \omega_s + c_2 \zeta_1 (D_m - \frac{4}{3} D_1) - \zeta_2 (c_4 D_m + c_6 D_2) + \zeta_m (c_7 D_m + 2c_7 D_2) \}_2 \\ + c_8 \{ 2\zeta_1 \bar{\omega}_1 + 2\zeta_2 \bar{\omega}_2 - (\zeta_m + 2\zeta_2) \bar{\omega}_3 \}_3 + c_8 \{ |k \cdot [-2V_1 \times \nabla \bar{\omega}_1 - 2V_2 \times \nabla \bar{\omega}_2 \\ + (V_m + 2V_2) \times \nabla \bar{\omega}_3] \}_4 \end{aligned} \quad (7.2.6)$$

The two remaining prognostic equations are computed from the difference between the vorticity equation for p_m and p_0 and the difference between the vorticity equation for p_1 and p_m . We will then have two vorticity equations valid for layer 1 and layer 2 respectively. The geostrophic terms are omitted.

$$\begin{aligned}
\frac{\partial \zeta_1}{\partial t} = & - \{ ((V_X)_m - (V_X)_1) \cdot \nabla \zeta_1 + (V_X)_1 \cdot \nabla (\eta_m - \zeta_1) \}_1 \\
& - \{ (\zeta_m - 2\zeta_1) D_1 + \zeta_1 D_m \}_2 + \left\{ \frac{\zeta_1}{p_0 - p_m} (\omega_m - \omega_s) \right\}_3 \\
& - \left\{ \frac{1}{(p_0 - p_m)} [k \cdot [(V_\psi)_1 \times \nabla (\omega_m - \omega_s)]] \right\}_4
\end{aligned} \tag{7.2.7}$$

$$\begin{aligned}
\frac{\partial \zeta_2}{\partial t} = & - \{ ((V_X)_m + (V_X)_2) \cdot \nabla \zeta_2 + (V_X)_2 \cdot \nabla (\eta_m + \zeta_2) \}_1 \\
& - \{ (\zeta_m + 2\zeta_2) D_2 + \zeta_2 D_m \}_2 + \left\{ \frac{1}{p_m - p_1} \zeta_2 (\omega_1 - \omega_m) \right\}_3 \\
& - \left\{ \frac{1}{p_m - p_1} [k \cdot [(V_\psi)_2 \times \nabla (\omega_1 - \omega_m)]] \right\}_4
\end{aligned} \tag{7.2.8}$$

Subscript 1 indicates contribution from $V_X \cdot \nabla \eta$

Subscript 2 indicates contribution from $\zeta \cdot \nabla V$

Subscript 3 indicates contribution from $\omega \frac{\partial \zeta}{\partial p}$

Subscript 4 indicates contribution from $k \cdot \left(\frac{\partial V}{\partial p} \times \nabla \omega \right)$

ω_m and ω_1 are the vertical motion at level p_m and p_1 respectively.

The non-geostrophic terms will be approximated by the divergence computed from the geostrophic part of the equations for the preceeding time step.

We can now express the forecasting equations in the following formal way (compare (2.2, 2.13 and 2.14).

$$\begin{aligned}
\nabla^2 \left\{ \frac{\partial}{\partial t} (\psi_m + c_1 \psi_2 - c_2 \psi_1) \right\} = & F_{mG} \\
+ F_{m1} + F_{m2} + F_{m3} + F_{m4}
\end{aligned} \tag{7.2.9}$$

$$\nabla^2 \frac{\partial \psi_1}{\partial t} - a_1 f^2 \frac{\partial \psi_1}{\partial t} + a_2 f^2 \frac{\partial \psi_2}{\partial t} = F_{1G} + F_{11} + F_{12} + F_{13} + F_{14} \quad (7.2.10)$$

$$\nabla^2 \frac{\partial \psi_2}{\partial t} - b_2 f^2 \frac{\partial \psi_2}{\partial t} + b_1 f^2 \frac{\partial \psi_1}{\partial t} = F_{2G} + F_{21} + F_{22} + F_{23} + F_{24} \quad (7.2.11)$$

F_{mG} , F_{1G} and F_{2G} are the geostrophic and non-adiabatic terms. They correspond to the right hand part of the equations 2.2, 2.13 and 2.14.

The non-geostrophic terms are the same as in the equations (7.2.6), (7.2.7) and (7.2.8).

8. NUMERICAL SOLUTION OF THE 3-PARAMETER MODEL

8.1 GENERAL ASPECTS

The equation of the model will be applied on a polar-stereographic projection. The polar-stereographic plane is assumed to cut the sphere at a latitude ϕ_0 . In the numerical computations ϕ_0 is put equal to 60°N , however, other values of ϕ_0 can easily be chosen. The grid-distance d can also be chosen arbitrarily. For further details see the program description in Appendix B.

The computational area can have any form, from an irregular octagon to a square. The only geometrical condition is that the inner angles of the area must be 90° or 135° .

The coordinate axis of the grid is positively oriented (see Figure 3). The computational area is specified by the coordinates of the corner points $(x_1/y_1 \dots x_8/y_8)$ and by the coordinates of the north pole $(x_{\text{pole}}/y_{\text{pole}})$.

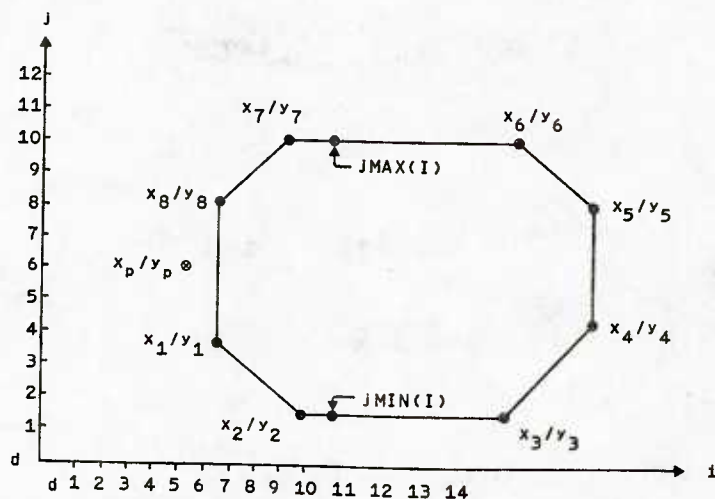


Figure 3. Example computational area.

If the area is reduced to a rectangle $x_1 = x_2$, $y_1 = y_2$,
 $x_3 = x_4$, $y_3 = y_4$, etc.

8.2 FINITE-DIFFERENCES

In order to transform the differential equations to finite-difference equations we will introduce the following finite-difference notation and operators (α and β are arbitrary quantities):

$$x \approx i\Delta s$$

$$y \approx j\Delta s$$

$$t \approx \tau\Delta t$$

$$f(x, y, t) \approx f_{i,j}^{\tau}$$

$$\nabla^2_{\alpha} \approx \frac{1}{d^2} \nabla^2_{\alpha} \quad (8.2.1a)$$

$$\mathcal{J}(\alpha, \beta) \approx \frac{1}{4d^2} \mathcal{J}(\alpha, \beta) \quad (8.2.1b)$$

$$\frac{\partial \alpha}{\partial t} \approx \frac{\Delta^{\tau} \alpha}{\epsilon \Delta t} \quad (8.2.1c)$$

$$(\nabla^2_{\alpha})_{ij} = \alpha_{i+1,j} + \alpha_{i-1,j} + \alpha_{i,j+1} + \alpha_{i,j-1} - 4\alpha_{ij} \quad (8.2.2a)$$

$$\begin{aligned} \mathcal{J}(\alpha; \beta) = & (\alpha_{i+1} - \alpha_{i-1})_j (\beta_{j+1} - \beta_{j-1})_i \\ & - (\alpha_{j+1} - \alpha_{j-1})_i (\beta_{i+1} - \beta_{i-1})_j \end{aligned} \quad (8.2.2b)$$

$$\tau = 0 \text{ yields } \alpha^{\frac{1}{2}} - \alpha^0 \text{ and } \epsilon = \frac{1}{2}$$

$$\tau = \frac{1}{2} \text{ yields } \alpha^1 - \alpha^0 \text{ and } \epsilon = 1$$

$$\tau \geq 1 \text{ yields } \alpha^{t+1} - \alpha^{t-1} \text{ and } \epsilon = 2 \quad (8.2.2c)$$

Introducing the map-scale factor $m = \frac{1+\sin\phi_0}{1+\sin\phi}$ and inserting the quantity $\mu = \left(\frac{m}{d}\right)^2$ we get:

$$\nabla^2(\Delta\psi_1) - a_1 \frac{f^2}{\mu} (\Delta\psi_2) + a_2 \frac{f^2}{\mu} (\Delta\psi_2) = F_{1G}, \quad (8.2.3)$$

$$\nabla^2(\Delta\psi_2) - b_2 \frac{f^2}{\mu} (\Delta\psi_2) + b_1 \frac{f^2}{\mu} (\Delta\psi_1) = F_{2G}, \quad (8.2.4)$$

$$\nabla^2(\Delta\psi_M) - \frac{q}{\mu} (\Delta\psi_M) = F_{mG}. \quad (8.2.5)$$

q is an empirical constant to adjust for the very long waves,
 $q = 0.75 \cdot 10^{-12} m^{-2}$.

Here we have (for simplicity we from now on put $J = J$):

$$F_{1G} = F_1' + F_1'',$$

$$F_{2G} = F_2' + F_2'',$$

$$F_1' = - \varepsilon \Delta t \cdot \frac{f}{\mu} (a_3 \omega_s + a_1 H),$$

$$F_1'' = - \varepsilon \Delta t \frac{1}{4} [J_1 + J_2 + f^2 (a_2 J_4 - a_1 J_3)],$$

$$F_2' = - \varepsilon \Delta t \frac{f}{\mu} (b_3 \omega_s + b_1 H),$$

$$F_2'' = - \varepsilon \Delta t \frac{1}{4} [J_5 + J_6 + f^2 (b_1 J_3 - b_2 J_4)],$$

$$F_{mG} = - \varepsilon \Delta t \frac{1}{4} [J_7 + J_8 + J_9 - 4 \frac{f}{\mu} c_8 \omega_s].$$

We further have:

$$\beta_2 = \eta_m^{-2} \zeta_1,$$

$$\beta_6 = \eta_m^{+2} \zeta_2,$$

$$\beta_7 = c_3 \eta_m^{-2} \zeta_1 + c_4 \zeta_2 + c_7 f,$$

$$\beta_8 = -c_2 \eta_m^{+2} \zeta_1,$$

$$\beta_9 = c_4 \eta_m^{+2} \zeta_2 + 2c_7 f,$$

$$\begin{aligned}
J_1 &= J(\psi_m; \zeta_1), \\
J_2 &= J(\psi_1; \beta_2), \\
J_3 &= J(\psi_m; \psi_1), \\
J_4 &= J(\psi_m; \psi_2), \\
J_5 &= J(\psi_m; \zeta_2), \\
J_6 &= J(\psi_2; \beta_6), \\
J_7 &= J(\psi_m; \beta_7), \\
J_8 &= J(\psi_1; \beta_8), \\
J_9 &= J(\psi_2; \beta_9), \\
J_{10} &= J(\psi_m - 2\psi_1; p_s), \\
\omega_s &= \frac{\mu}{4} J_{10} - c_f (\zeta_m - 2\zeta_1).
\end{aligned}$$

c_f is an empirical constant and a function of the exchange coefficient of momentum in the boundary layer.

8.3 COMPUTATION OF SENSIBLE HEAT

Equation (4.1.4) reads in finite-difference form:

$$\begin{aligned}
H_{\text{SENS}} &= 0.5 \cdot 10^{-2} h (0.1) \sqrt{\frac{\mu}{4} ((\Delta_x \psi_o)^2 + (\Delta_y \psi_o)^2) + 1} (T_s - h_2 \psi_1) \\
&\quad \text{over ocean, if } (T_s - h_2 \psi_1) > 0, \quad (8.3.1)
\end{aligned}$$

$$H_{\text{SENS}} = 0 \text{ over land and if } (T_s - h_2 \psi_1) < 0.$$

$$\Delta_x \psi_o = [\psi_m(i+1) - 2\psi_1(i+1)] - [\psi_m(i-1) - 2\psi_1(i-1)]$$

$$\Delta_y \psi_o = [\psi_m(j+1) - 2\psi_1(j+1)] - [\psi_m(j-1) - 2\psi_1(j-1)]$$

8.4 COMPUTATION OF LATENT HEAT (See Gambo 1963)

The latent heat is introduced in the model by the expression:

$$\begin{aligned} H_{LAT} &= -hL\omega^*F^* && \text{if } \bar{\omega}_1 \leq -\delta_1 \\ H_{LAT} &= 0 && \text{if } \bar{\omega}_1 > -\delta_1 \end{aligned} \quad (8.4.1)$$

$$H_{LAT} = 0 \quad \text{if } \frac{r}{\epsilon \Delta t} < \text{tolerance}$$

$$\begin{aligned} \omega^* &= -(\delta_1 + \delta_2) \quad \text{if } \bar{\omega}_1 < -(\delta_1 + \delta_2) \\ \omega^* &= -\left| \frac{\bar{\omega}_1^2}{\delta_1 + \delta_2} \right| \quad \text{if } -(\delta_1 + \delta_2) \leq \bar{\omega}_1 \leq -\delta_1 \end{aligned} \quad (8.4.2)$$

δ_1 and δ_2 are here two tolerances with the same dimension as $\bar{\omega}_1$, r is the precipitation, and δ_2 is introduced for operational purposes.

$$F^* = F^*(p, T) = \frac{\epsilon \frac{T}{p} E(T) \left[\frac{\epsilon L}{C_p} - T \right]}{pT^2 + \frac{\epsilon^2 L^2}{C_p R} E(T)} \quad (8.4.3)$$

$$E(T) = E_0 e^{\frac{\epsilon L}{R} \left(\frac{1}{T_0} - \frac{1}{T} \right)}$$

$$\epsilon = 0.622$$

$$L = 2.5 \cdot 10^6$$

$$P = P_{\text{mean}} = \frac{P_m + P_0}{2}$$

$$C_p = 1004$$

$$T_0 = 273$$

$$E_0 = 0.611$$

$$T = \frac{2f_0}{P \ln\left(\frac{P_0}{P_m}\right)} \psi_1 = h_6 \psi_1$$

8.4.1 Computation of \tilde{W} and \tilde{D}

We will compute the humidity in the layer

$P_O = 100$ cb and $P_T = 30$ cb; defining \tilde{W} as

$$\tilde{W} = \frac{1}{70} \int_{30}^{100} W(P) E(P) dp$$

and inserting the wind profile yields

$$\tilde{W} = e_m W_m + e_1 W_1 + e_2 W_2 \quad (8.4.4)$$

where

$$\begin{aligned} e_m &= \frac{1}{70} [15E(100) a_{100} + 25E(70) a_{70} + 20E(50) a_{50} + 10E(30) a_{30}] , \\ e_1 &= \frac{1}{70} [15E(100) b_{100} + 25E(70) b_{70} + 20E(50) b_{50} + 10E(30) b_{30}] , \\ e_2 &= \frac{1}{70} [15E(100) c_{100} + 25E(70) c_{70} + 20E(50) c_{50} + 10E(30) c_{30}] . \end{aligned} \quad (8.4.5)$$

The constants a_p , b_p and c_p , where p assumes the values 100, 70, 50, 30 are computed for the three following alternatives:

I	II	III
$p_m \leq p \leq p_O$	$p_1 \leq p \leq p_m$	$0 \leq p \leq p_1$
$a_p = 1$	$a_p = 1$	$a_p = \frac{p}{p_1}$
$b_p = - \frac{2(p-p_m)}{p_O-p_m}$	$b_p = 0$	$b_p = 0$
$c_p = 0$	$c_p = \frac{2(p_m-p)}{p_m-p_1}$	$c_p = \frac{2p}{p_1}$

$$E(100)=2.5415, E(70)=0.9683, E(50)=0.3527, E(30)=0.0613$$

Correspondingly we get

$$\tilde{D} = e_m D_m + e_1 D_1 + e_2 D_2 = (e_1 + e_m c_2) D_1 + (e_2 - e_m c_1) D_2 - e_m c_8 w_s$$

8.4.2 Computation of Precipitation

Equation (4.2.6) reads in finite-difference form (the term $\nabla \tilde{\chi} \nabla w$ disregarded and w replaced with q):

$$q^{\tau+1} = q^{\tau-1} + \varepsilon \Delta t H_q^\tau,$$

$$H_q^\tau = -\frac{\mu}{4} J(\tilde{\psi}^\tau; q^\tau) - q^\tau (\tilde{D}^\tau + dw_s^\tau) + A_{diff} \mu \nabla^2 q^\tau. \quad (8.4.6)$$

The precipitation is computed in the following way (precipitation is indicated by r):

$$q_{SAT} = q(\bar{T}_\psi) = q(h_3 \psi_1 + h_4 \psi_2)$$

$$\text{If } (q^{\tau+1} - 0.8 q_{SAT}) > 0 \text{ then } q^{\tau+1} = 0.8 q_{SAT},$$

$$r = \Delta p (q^{\tau+1} - 0.8 q_{SAT}).$$

$$\text{If } (q^{\tau+1} - 0.8 q_{SAT}) < 0 \text{ then } r = 0 \text{ and if}$$

$$(q^{\tau+1} - 0.2 q_{SAT}) < 0 \text{ then } q_{mod}^{\tau+1} = 0.2 q_{SAT}. \quad (8.4.7)$$

r is accumulated for every timestep and printed out at certain prescribed times. See Appendix B.

8.5 NUMERICAL SOLUTION OF THE NON-GEOSTROPHIC TERMS

The finite difference analogues for (7.2.9), (7.2.10) and (7.2.11) read:

$$\nabla^2 (\Delta\psi_M) - \frac{g}{\mu} (\Delta\psi_M) = F_{mG} + \frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{mi} \quad (8.5.1)$$

$$\nabla^2 (\Delta\psi_1) - a_1 \frac{f^2}{\mu} (\Delta\psi_1) + a_2 \frac{f^2}{\mu} (\Delta\psi_2) = F_{1G} + \frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{1i} \quad (8.5.2)$$

$$\nabla^2 (\Delta\psi_2) - b_2 \frac{f^2}{\mu} (\Delta\psi_2) + b_1 \frac{f^2}{\mu} (\Delta\psi_1) = F_{2G} + \frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{2i} \quad (8.5.3)$$

$$\begin{aligned} F_{m1} = & - \frac{\mu}{2} \{ \nabla \chi_m \cdot \nabla [\underbrace{c_3(\zeta_m + f) - c_2\zeta_1 + c_4\zeta_2 + c_7f}_{\beta_7}] \\ & + \nabla \chi_1 \cdot \nabla [\underbrace{-c_2(\zeta_m + f) + c_5\zeta_1}_{\beta_8}] \\ & + \nabla \chi_2 \cdot \nabla [\underbrace{c_4(\zeta_m + f) + c_6\zeta_2 + 2c_7f}_{\beta_9}] \} \end{aligned} \quad (8.5.4a)$$

$$\begin{aligned} F_{m2} = & [D_1(k_1\zeta_1 + k_2\zeta_2 + k_3\zeta_m) + D_2(k_4\zeta_1 + k_5\zeta_2 + k_6\zeta_m) \\ & + \omega_s(k_7\zeta_1 + k_8\zeta_2 + k_9\zeta_m)] \end{aligned} \quad (8.5.4b)$$

$$\begin{aligned} F_{m3} = & [D_1(k_{10}\zeta_1 + k_{11}\zeta_2 + k_{12}\zeta_m) + D_2(k_{13}\zeta_1 + k_{14}\zeta_2 + k_{15}\zeta_m) \\ & + \omega_s(k_{16}\zeta_1 + k_{17}\zeta_2 + k_{18}\zeta_m)] \end{aligned} \quad (8.5.4c)$$

$$F_{m4} = \frac{\mu}{2} \{ \nabla \psi_1 \cdot \nabla [\underbrace{k_{19}D_1 + k_{20}D_2 + k_{21}\omega_s}_{\gamma_7}] \}$$

$$\begin{aligned}
& + \nabla \psi_2 \cdot \nabla \underbrace{[k_{22} D_1 + k_{23} D_2 + k_{24} \omega_s]}_{\gamma_8} \\
& + \nabla \psi_m \cdot \nabla \underbrace{[k_{25} D_1 + k_{26} D_2 + k_{27} \omega_s]}_{\gamma_9} \} \quad (8.5.4d)
\end{aligned}$$

$$F_{11} = - \frac{\mu}{2} \{ \nabla \chi_m \cdot \nabla \zeta_1 + \nabla \chi_1 \cdot \nabla (\zeta_m + f - 2\zeta_1) \} \quad (8.5.5a)$$

$$F_{12} = [D_1 (k_{28} \zeta_1 - \zeta_m) + D_2 k_{29} \zeta_1 + \omega_s k_{30} \zeta_1] \quad (8.5.5b)$$

$$F_{13} = [\zeta_1 (k_{31} D_1 + k_{32} D_2 + k_{33} \omega_s)] \quad (8.5.5c)$$

$$\begin{aligned}
F_{14} = \frac{\mu}{2} \{ \nabla \psi_1 \cdot \nabla \underbrace{[k_{31} D_1 + k_{32} D_2 + k_{33} \omega_s]}_{\gamma_{10}} \} \\
\quad (8.5.5d)
\end{aligned}$$

$$F_{21} = - \frac{\mu}{2} \{ \nabla \chi_m \cdot \nabla \zeta_2 + \nabla \chi_2 \cdot \nabla (\zeta_m + f + 2\zeta_2) \} \quad (8.5.6a)$$

$$F_{22} = \{ D_1 k_{34} \zeta_2 + D_2 [k_{35} \zeta_2 - \zeta_m] + \omega_s k_{30} \zeta_2 \} \quad (8.5.6b)$$

$$F_{23} = [\zeta_2 (k_{36} D_1 + k_{37} D_2 + k_{38} \omega_s)] \quad (8.5.6c)$$

$$\begin{aligned}
F_{24} = \frac{\mu}{2} \{ \nabla \psi_2 \cdot \nabla \underbrace{[k_{36} D_1 + k_{37} D_2 + k_{38} \omega_s]}_{\gamma_{11}} \} \\
\quad (8.5.6d)
\end{aligned}$$

The computation of the non-geostrophic forcing function

$$\frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{mi}, \quad \frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{li} \quad \text{and} \quad \frac{\varepsilon \Delta t}{\mu} \sum_{i=1}^4 F_{2i}$$

are performed by a separate program. See Appendix B.

8.5.1 Numerical Coefficients for the Non-geostrophic Terms

Expressions for constants of the non-geostrophic terms
read:

$$\bar{\omega}_3 = t_6 \omega_s + t_7 D_1 + t_8 D_2$$

$$\omega_m = t_9 \omega_s + t_{10} D_1 + t_{11} D_2$$

$$\omega_1 = t_{12} \omega_s + t_{13} D_1 + t_{14} D_2$$

$$t_6 = c_8 \frac{p_1}{6}$$

$$t_7 = -c_2 \frac{p_1}{6}$$

$$t_8 = (c_1 - 2) \frac{p_1}{6}$$

$$t_9 = 1 - c_8 (p_o - p_m)$$

$$t_{10} = (c_2 - 1) (p_o - p_m)$$

$$t_{11} = -c_1 (p_o - p_m)$$

$$t_{12} = c_8 \frac{p_1}{2}$$

$$t_{13} = -c_2 \frac{p_1}{2}$$

$$t_{14} = (c_1 - 2) \frac{p_1}{2}$$

$$k_1 = c_2 (c_2 - \frac{4}{3})$$

$$k_2 = -c_2 c_4$$

$$k_3 = c_2 c_7$$

$$k_4 = -c_1 c_2$$

$$k_5 = c_1 c_4 - c_6$$

$$k_6 = c_7 (2 - c_1)$$

$$k_7 = -c_2 c_8$$

$$k_8 = c_4 c_8$$

$$k_9 = c_8 (1 - c_7)$$

$$k_{10} = 2t_2 c_8$$

$$k_{11} = -2(t_3 + t_7) c_8$$

$$k_{12} = -t_7 c_8$$

$$k_{13} = -2t_3 c_8$$

$$k_{14} = 2(t_5 - t_8) c_8$$

$$k_{15} = -t_8 c_8$$

$$k_{16} = 2t_1 c_8$$

$$k_{17} = 2(t_4 - t_6) c_8$$

$$k_{18} = -t_6 c_8$$

$$k_{19} = 2t_2 c_8$$

$$k_{20} = -2t_3 c_8$$

$$k_{21} = 2t_1c_8$$

$$k_{22} = -2(t_3+t_7)c_8$$

$$k_{23} = 2(t_5-t_8)c_8$$

$$k_{24} = 2(t_4-t_6)c_8$$

$$k_{25} = -t_7c_8$$

$$k_{26} = -t_8c_8$$

$$k_{27} = -t_6c_8$$

$$k_{28} = 2-c_2$$

$$k_{29} = +c_1$$

$$k_{30} = +c_8$$

$$k_{31} = -\frac{t_{10}}{(p_o-p_m)}$$

$$k_{32} = -\frac{t_{11}}{(p_o-p_m)}$$

$$k_{33} = \frac{1-t_9}{(p_o-p_m)}$$

$$k_{34} = -c_2$$

$$k_{35} = c_1-2$$

$$k_{36} = \frac{t_{10}-t_{13}}{(p_m-p_1)}$$

$$k_{37} = \frac{t_{11} - t_{14}}{(p_m - p_1)}$$

$$k_{38} = \frac{t_9 - t_{12}}{(p_m - p_1)}$$

For $p_o=100$, $p_m=50$ and $p_1=30$ we have the following numerical values for the constants:

$$t_6 = 0.0588235$$

$$t_7 = -2.941175$$

$$t_8 = -7.058825$$

$$t_9 = 0.411765$$

$$t_{10} = -20.588250$$

$$t_{11} = -29.411750$$

$$t_{12} = 0.176471$$

$$t_{13} = -8.823525$$

$$t_{14} = -21.176475$$

$$k_1 = -0.438291$$

$$k_2 = -0.276816$$

$$k_3 = 0.034602$$

$$k_4 = -0.346020$$

$$k_5 = -0.507498$$

$$k_6 = 0.083045$$

$$k_7 = -0.006920$$

$$k_8 = 0.005536$$

$$k_9 = 0.011073$$

$$k_{10} = -0.438294$$

$$k_{11} = -0.276817$$

$$\begin{aligned}
k_{12} &= 0.034602 \\
k_{13} &= -0.346021 \\
k_{14} &= -0.507498 \\
k_{15} &= 0.083045 \\
k_{16} &= 0.016609 \\
k_{17} &= 0.005536 \\
k_{18} &= -0.000692 \\
k_{19} &= -0.438293 \\
k_{20} &= -0.346021 \\
k_{21} &= 0.016609 \\
k_{22} &= -0.276817 \\
k_{23} &= -0.507497 \\
k_{24} &= 0.005536 \\
k_{25} &= 0.034602 \\
k_{26} &= 0.083045 \\
k_{27} &= -0.000692 \\
k_{28} &= 1.411765 \\
k_{29} &= +0.588235 \\
k_{30} &= +0.0117647 \\
k_{31} &= 0.411765 \\
k_{32} &= 0.588235 \\
k_{33} &= 0.0117647 \\
k_{34} &= -0.588235 \\
k_{35} &= -1.411765
\end{aligned}$$

$$k_{36} = -0.588235$$

$$k_{37} = -0.411765$$

$$k_{38} = 0.0117647$$

9. INITIALIZATION

The stream functions ψ are computed from the geopotential Z by the relation

$$\begin{aligned} g\nabla^2 Z &= \nabla \cdot (f\nabla\psi) \quad \text{or} \\ \nabla^2 \psi &= \frac{g}{f}\nabla^2 Z - \frac{1}{f}\nabla f \cdot \nabla\psi; \end{aligned} \quad (9.1)$$

inserting

$$\begin{aligned} \nabla\psi &= \frac{g}{f}\nabla Z \quad \text{yields} \\ \nabla^2 \psi &= \frac{g}{f}\nabla^2 Z - \frac{g}{f^2}\nabla f \cdot \nabla Z. \end{aligned} \quad (9.2)$$

The boundary values of ψ are computed by the relation (γ is a constant)

$$\frac{\partial\psi}{\partial s} = \frac{g}{f} \frac{\partial Z}{\partial s} + \gamma. \quad (9.3)$$

If we assume that there is no net flow out of the area it is easily seen that $\gamma = -\frac{1}{L} \oint \frac{g}{f} \frac{\partial Z}{\partial s} ds$. The integration is performed along the boundary of the area in positive order. The integration starts in point x_1/y_1 where we put

$$\psi(x_1; y_1) = \frac{g}{f(x_1; y_1)} Z(x_1; y_1). \quad (9.4)$$

Z is computed from ψ in a similar way using the following equation:

$$\nabla^2 Z = \frac{f}{g}\nabla^2 \psi + \frac{1}{g}\nabla f \cdot \nabla\psi. \quad (9.5)$$

Boundary values for Z are computed initially and stored in a special string. See Appendix B.

9.1 NUMERICAL SOLUTION

Equation (9.2) reads in finite difference form:

$$\nabla^2 \psi = \frac{g}{f} \nabla^2 Z - \frac{g}{2f^2} \nabla f \cdot \nabla Z \quad (9.6)$$

where the standard five point formulas are used to compute $\nabla^2 \psi$, $\nabla^2 Z$ and $\nabla f \cdot \nabla Z$ is defined as

$$\begin{aligned} \nabla f \cdot \nabla Z = & (f_{i+1} - f_i)_j (Z_{i+1} - Z_i)_j + (f_i - f_{i-1})_j (Z_i - Z_{i-1})_j \\ & + (f_{i+1} - f_j)_i (Z_{i+1} - Z_j)_i + (f_j - f_{j-1})_i (Z_j - Z_{j-1})_i. \end{aligned}$$

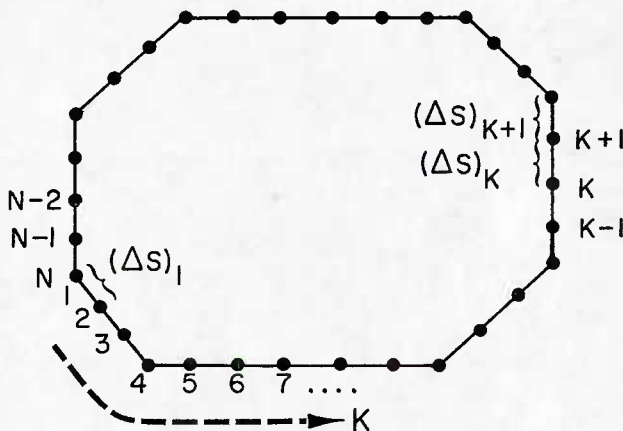
Equation (9.3) reads in finite difference form:

$$\psi_{k+1} = \psi_k + 2g \frac{Z_{k+1} - Z_k}{f_{k+1} + f_k} + \gamma (\Delta s)_k \quad (9.7)$$

where

$$\gamma = - \frac{1}{L} 2g \sum_{k=1}^{N-1} \frac{Z_{k+1} - Z_k}{f_{k+1} + f_k} \quad (9.8)$$

Figure 4 defines k and the order of integration.



Z to ψ and ψ to
 Z is computed by
the program STREAMF
(see Appendix B).

Figure 4. Definition of k and order of integration

9.2 INITIALIZATION OF THE SPECIFIC HUMIDITY

Since we are using the ψ -functions as the time-dependent variables, we have to modify the initial humidities in order to make them consistent with the time-integration.

Therefore, we put

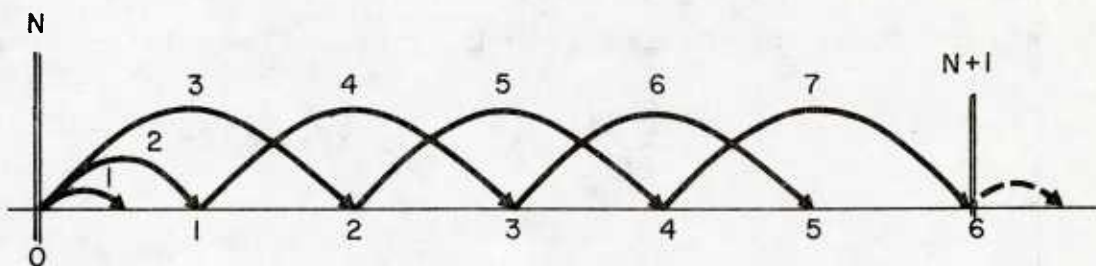
$$q = q_{\text{analyzed}} \cdot \frac{q_{\text{sat}}(\bar{T}_{\psi})}{q_{\text{sat}}(\bar{T}_Z)} \quad (9.2.1)$$

$$\bar{T}_Z = \frac{2g}{R \ln 2} \left(\frac{Z_{500} - Z_{1000}}{2} \right) = \frac{g}{f_0} (h_3 Z_1 + h_4 Z_2), \text{ and}$$

$$\bar{T}_{\psi} = h_3 \psi_1 + h_4 \psi_2.$$

10. TIME-INTEGRATION AND TREATMENT OF LATERAL BOUNDARY VALUES

The forecast is basically computed in 6-hour intervals but this interval can easily be changed. The first time step in each interval is non-centered. Smoothing, elliptization and printing of results (if so desired) are performed at the end of every interval.



$$\frac{\partial \psi}{\partial t} \sim \frac{\Delta \psi}{\epsilon \Delta t} \quad \Delta t = 1 \text{ hour in this example} \quad (10.1)$$

ND number of Δt for the interval

kT time step index; $k = 1, 2, 3 \dots ND+1$

N index for every interval integration (e.g., 6-hour interval)

N	Forecast length
0	0
1	+6
2	+12
3	+18
.	.
.	.
.	.

Initial height fields Z^0 are stored on secondary storage during the whole computation. In the case of variable boundary conditions, Z^0 is followed by Z^N ($N=1,2,3\dots$) (forecasts for each interval).

Assuming

$$\frac{\partial \psi}{\partial t} = \frac{g}{f} \frac{\partial Z}{\partial t} \quad \text{or} \quad \psi^N - \psi^0 = \frac{g}{f} (Z^N - Z^0) \quad (10.2)$$

the stream function at time step k can be interpolated from:

$$\psi^k = (\psi^0 - \frac{g}{f} Z^0) + \frac{g}{f} [(1-\alpha_k) Z^N + \alpha_k Z^{N+1}] \quad (10.3)$$

where

$$\alpha_k = \frac{kT-1}{ND}$$

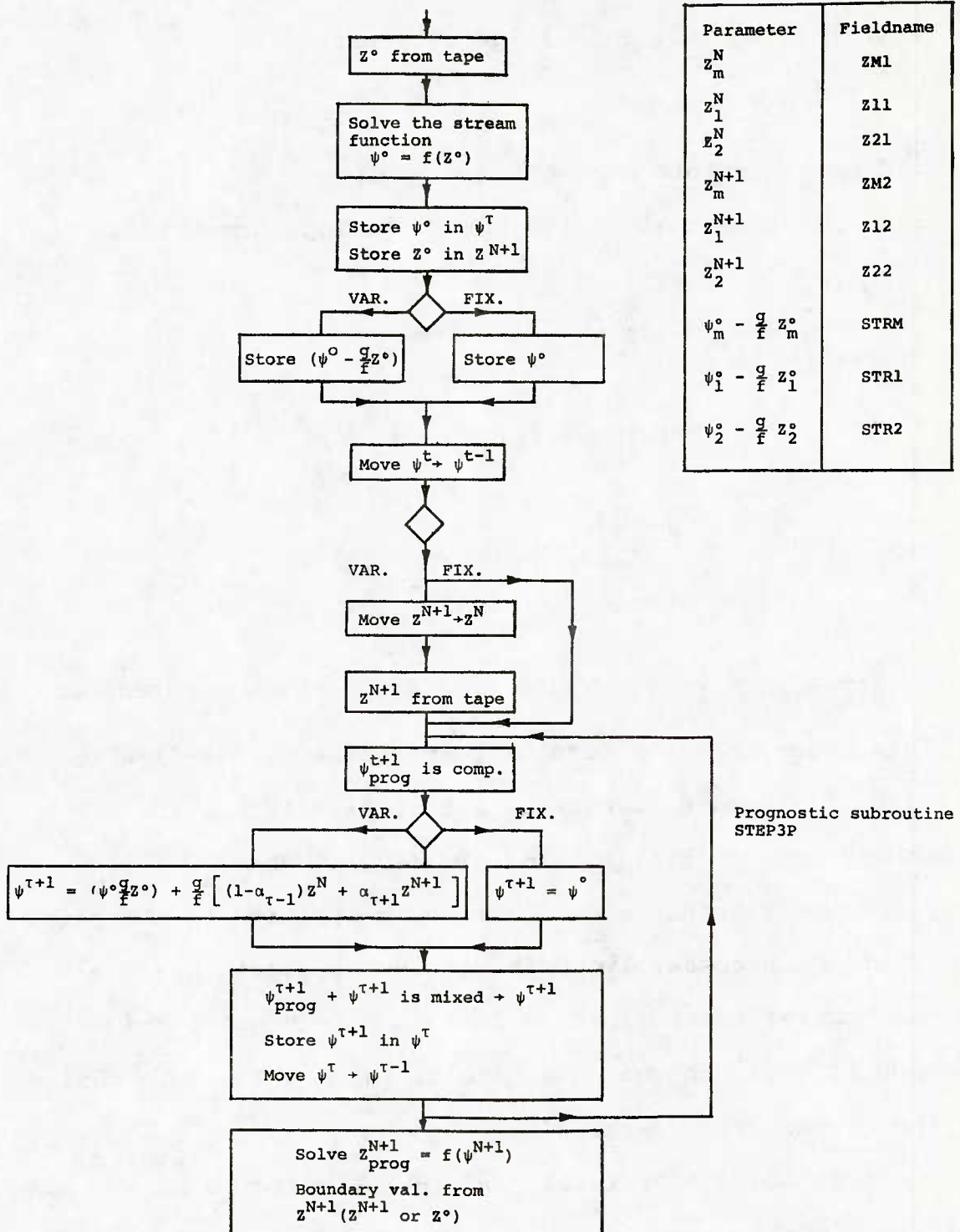
At each time step this interpolated stream function is mixed with the forecasted stream function ψ_{prog}^k in the following way:

Boundary values	$\psi_{\text{mod}} = \psi_{\text{prog}}^k + w1(\psi^k - \psi_{\text{prog}}^k)$	
1 grid point inside the boundary	$\psi_{\text{mod}} = \psi_{\text{prog}}^k + w2(\psi^k - \psi_{\text{prog}}^k)$	
2 grid points inside the boundary	$\psi_{\text{mod}} = \psi_{\text{prog}}^k + w3(\psi^k - \psi_{\text{prog}}^k)$	
3 grid points or more inside the boundary	$\psi_{\text{mod}} = \psi_{\text{prog}}^k$	(10.4)

$w1, w2$ and $w3$ are 3 predetermined constants with $0 \leq w1, w2, w3 \leq 1$

The following flow diagram illustrates the treatment of the boundary values:

Flow diagram: Treatment of boundary values.



11. THE CRITERION OF ELLIPTICITY

An iterative procedure is used to modify a stream-function field so that the ellipticity criterion

$$\nabla^2 \psi + \frac{f}{2} > 0$$

is valid in all points of the field.

Each point is tested with the following formula:

$$\mu \nabla^2 \psi + \frac{f}{2} - \epsilon = \delta$$

where

$$\nabla^2 \psi = \psi_{i+1,j} + \psi_{i-1,j} + \psi_{i,j+1} + \psi_{i,j-1} - 4 \cdot \psi_{i,j}$$

and

$$\epsilon = 0.001 \cdot f.$$

If $\delta < 0$, $\psi_{i,j}$ is modified by

$$\psi_{i,j} = \psi_{i,j} - \frac{\epsilon - \delta}{2\mu} \cdot k \text{ where } k \text{ is a convergence parameter.}$$

This means that the vorticity increases by $2(\epsilon - \delta) \cdot k$ in the point (i,j) and decreases by $0.5(\epsilon - \delta) \cdot k$ in the four surrounding points $(i+1,j)$, $(i-1,j)$, $(i,j+1)$ and $(i,j-1)$. This procedure guarantees that δ becomes positive in the point (i,j) , but not necessarily in the surrounding points. The test must therefore be repeated for all points until the criterion is valid in the whole field. With a suitable choice of k the method is convergent.

k can be found empirically and is estimated to be of the order $k \approx 0.85$.

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APPENDIXES

APPENDIX A
GENERAL PRINCIPLES
IN THE
PROGRAMMING OF THE THREE PARAMETER MODEL

The programming has been performed in a very general way. It is therefore possible:

- a. to generate initial fields and run the model in a channel with variable cyclic boundary conditions (this is the case for the actual program).
- b. to run the model with constant boundary values on a given area (see above).
- c. to run the model with variable boundary values, the boundary values taken from another model.

The model needs 13 fields in the core;

1 field for the coriolis parameter	F,
1 field for $\mu = \frac{m^2}{d^2}$	MY,
1 special field indicator	MARK,
10 fields for the model computations	F1-F10.

In addition to this, 40 fields (28 for the quasi-geostrophic part of the model) are specified on secondary memories (disks, drums or large extended core).

"Household" programs:

PUT1: This program computes FPAR (see common,
 Subroutine JMIMA)

MARKF: This program computes a special integer field MARK. The program MARK specifies status of the field. Points outside the area = 0, points inside the are <0 and points on the boundary >0. The corner points etc. are specified according to given examples.

MYFF: This program computes f and μ and puts the result in the fields F and MY.

RANWT: Writes fields on secondary storage e.g., disk drums, or extended core storage.

RANRD: Reads fields from secondary storage e.g., disk, drums, or extended core storage.

MAP: Prints pattern on line printer; 0 (zero) points, resolution and map scale are specified.

GENCH: Generates initial state for a channel flow.

BMOVE: Administrates computation of cyclic boundary conditions for a channel flow.

Since the axis of the grid is defined differently from what is used in Fortran, the programming of finite difference operators should be performed in a special way. This is not necessary, but speeds up the computation considerably. As an example, subroutine JACOB reads:

.
.
.
.
.

MI = M-1

Dφ 10 I = 2,MI

J1 = JMIN(I)+1

J2 = JMAX(I)-1

K = (J1-1)*M+I

Dφ 10 J = J1,J2

C(K) = (A(K+1) - A(K-1)*(B(K+M) - B(K-M)).....

10 K = K + M

.
.
.

APPENDIX B
PROGRAM SPECIFICATIONS

The following programs are written for the 3-parameter model.

Level 1: Main program

Level 2: Subroutines called by level 1

Level 3: Subroutines called by level 2

Level 4: Subroutines called by level 3

<u>Level</u>	<u>Program</u>
1	PROG3P
2	PUT1
3	JMIMA
2	COEFF3P
2	MARKF
2	MYFF
2	STREAMF
3	BDRGDR
3	BDRVAL
2	SATUR
2	RANWT (RANRD)
2	ELLIPT
2	GENCH
2	ASMUT
2	MAP3P
3	MAP
2	STEP3P

<u>Level</u>	<u>Program</u>
3	JACOB
3	ABSVOR (RELVOR)
3	MIXF
3	HELM (POIS)
3	HELMSYS
3	BMOVE
3	STEPEXT
4	GRADPR
4	VELPOT

Program PROG3 (Main program for 3-parameter model, see flow diagram B-1 and description.)

Arrays

<u>Name</u>	<u>Dimensions</u>	<u>Description</u>
F1-F10	2	Working fields in core. See flow diagram B-1.
F,MY (real)	2	Fields for Coriolis parameter, f , and $\mu = (\frac{m}{d})^2$. m is the mapscale factor for a polar stereographic projection and d is the grid length. F and MY are generated by the subroutine MYFF.
MARK	2	Marking of each grid point by a special integer (index). See description. MARK is generated by the subroutine MARKF.
UPS, UPM, UP1	1	Zonal wind profiles at the levels p_s , p_m and p_1 for the initial fields in the channel case. The values are introduced by DATA statements.
WX	1	Working array for generation of initial fields in the channel case. Dimension shall be at least equal to the number of grid-points across the channel.

<u>Name</u>	<u>Dimensions</u>	<u>Description</u>
NX,NY	1	Wave numbers for perturbations of the initial fields in the channel case in x and y directions. The values are introduced by DATA statements.
PSIC, PSIS, LAMC, LAMS	1	Amplitudes (PSI) and phase lags (LAM) perturbations of the initial field with cosine and sine profiles in the y direction respectively. For both profiles a sine wave is used in the x-direction. The values are introduced by DATA statements. See further description of GENCH.
ZB,PSIB,FB,SB	1	Boundary strings for storage of boundary values in counterclockwise order, starting with corner point 1 (see FPAR). Used by the subroutine STREAMF.
IT	2	Storage of the number of iterations for the solution of a): the system of two Helmholtz equations and b): the Helmholtz equation for the mean field. Values for an integration interval are stored in the array by the subroutine STEP3P.

<u>Name</u>	<u>Dimensions</u>	<u>Description</u>
MAPHOUR (1)	1	Integers giving the hours for mapping (2).
NSMUTT (1)	1	Integers giving the hours for smoothing (2).
MELLIPT (1)	1	Integers giving the hours for ellipticity (2).
LETACC (1)	1	Integers giving the hours at which the accumulation of precipitation shall be interrupted (2).

(1) The values shall be given in DATA statements and must be multiples of the integration interval (e.g., 6 hours). Unused positions are indicated by -1.

(2) This value can be increased if this is necessary.

Flow Diagram B-1: PROG3P

F1	F2	F3	F4	F5	F6	F7	F8	F9	F10
Z ^O (P _S)	Z ^O (P _m)	Z ^O (P _L)	q ^{ONA}	P _S	T _S	ψ ^O (P _S)	ψ ^O (P _m)	ψ ^O (P _L)	F10
Z ^O _m	Z ^O _L	Z ^O ₂	q			x	x	x	
+	+	+				ψ ^O _m	ψ ^O _L	ψ ^O ₂	
ψ ^O _m	ψ ^O _L	ψ ^O ₂				ψ ^O _m -g ^O ₂ / _m	ψ ^O _L -g ^O ₂ / _L	ψ ^O ₂ -g ^O ₂ / ₂	
ψ ^O _m	ψ ^O _L	ψ ^O ₂				ψ ^O _m	ψ ^O _L	ψ ^O ₂	
ψ ^O _m	ψ ^O _L	ψ ^O ₂				+	+	+	
ψ ^O _m	ψ ^O _L	ψ ^O ₂	q ^O (~50%)	(100.0)	(273.0)	ψ ^O (P _S)	ψ ^O (P _m)	ψ ^O (P _L)	
+	+	+	+	+	+				
+	+	+							0.0
0.0	x								
Z ^{T+L} _Z (P _S)	Z ^{T+L} _Z (P _m)	Z ^{T+L} _Z (P _L)							
Z ^{T+L} _m	Z ^{T+L} _L	Z ^{T+L} ₂							
+	+	+							

INITF
STREAMF

→ZM2,Z12,Z22
ELLIPT

var. bound.

→STRM,STR1,STR2

GENCH

→STRM,STR1,STR2

→PSIM1,PSI11,PSI21,
HUM1,PS,TS,
O→PREC,WS,DIV1,DIV2

O→DIV1,DIV2
HUM1→HUM2,PSIM1→PSIM2
PSI11→PSI12,PSI12→PSI22
ZM2→ZM1,Z12→Z11
Z22→Z21

ZINPUT

→ZM2,Z12,Z22

STEP3P

SMOOTHING

ELLIPT

O→PREC

Common areas

/FPAR/

<u>Name</u>	<u>Type</u>	<u>Description</u>
IC(8),JC(8)	integer	i- and j-coordinates for the corner points of the area.
XPOL,YPOL	real	i- and j-coordinates for the north pole. Specified relative to the origin of the grid.
R	real	radius of the earth.
RE	real	Distance from the pole to the equator on the map in grid-length units.
DS	real	grid-length (in meters).
JMIN(100), JMAX(100)	integer	minimum and maximum j-coordinate for each i-column in the field.
M,N	integer	Dimension of the field arrays.
KIND	integer	Channel indicator. Channel = 1, No channel = 0. Value shall be given by DATA statement.

/FORM1/

IC1(8),JC1(8)	integer	Same as IC(8),JC(8) for actual field. Values introduced by DATA statement.
---------------	---------	--

<u>Name</u>	<u>Type</u>	<u>Description</u>
JMIN1(100), JMAX1(100)	integer	Same as JMIN(100), JMAX(100)
XP,YP,D1	real	Same as XPOL, YPOL, DS. Values shall be introduced by DATA statements. D1 is given in km.
/KANAL/		
FIM	real	Used for channel computations. Latitude for the middle of the channel. Used for computation of B-plane. To be given by DATA statement.
/DRM/*		
MN	integer	=M*N. To be computed in main program
NDIM	integer	Core memory space available for field arrays. The arrays F1-F10, F,MY,MARK must be included in this area. Additional space is used for storage of fields in COMMON area//.
NFLD	integer	Number of field arrays in core
SL	real	Working area for core memory fields. The dimension must be at least SL(NDIM)

Arrays stored at extended core, disks or drums. 1 in the end
of the name indicate timestep τ , 2 in the end of the name
timestep $\tau-1$.

* Not necessary to specify if only extended core storage is used.

PSIM1,PSI11,PSI21 PSIM2,PSI12,PSI22	real	Storage arrays for stream functions. See flow diagram B-1.
HUM1,HUM2,DIV1 DIV2,WS,HEAT	real	Storage arrays for humidity, divergence, vertical velocity at the lower boundary, and heat.
J789,J12,J56,J3		Storage arrays for Jacobians.
PS,TS,PREC		Storage arrays for surface standard pressure, sea surface temperature, and accumulated precipitation.
STRM,STR1,STR2,ZM1 Z11,Z21,ZM2,Z12,Z22	real	Storage arrays. See Chapter 10 (lateral boundary mixing).
H11,H21,HM1		Fields for advection of vorticity by the divergent wind, thickness field (1,2) and mean field (M).
H12,H22,HM2		$\omega \frac{\partial \zeta}{\partial t} + \zeta \nabla \cdot \mathbf{V}$ for thickness fields (1,2) and mean field (M).
H13,H23,HM3		Twisting term for thickness fields (1,2) and for mean field, M.
V1,V2,VM		Velocity potential for thickness fields (1,2) and for mean field (M).
ID(200)*	integer	Catalog array for direct memory access.

* Not necessary to specify if only extended core storage is used.

/COEFF/

A1,A2.....,T5	real	Constants used in the model. Computed by the subroutine COEFF3P from given values on stability and pressure levels.
---------------	------	--

/COEFF2/

T6,TF.....,T38	real	Constants used for the computation of the non-geostrophic terms. Computed by COEFF3P from given values on stability and pressure levels.
----------------	------	--

/RUNPAR/

DELT	real	Timestep in sec. computed in the main program.
NTSTEP	integer	Number of timesteps for an inte- gration interval (6 hours). To be defined in a DATA statement.
ALFASYS, ALFAM ALFAZ, ALFAPSI, RESSYS, RESM, RESZ, RESPSI	real	Overrelaxation coefficients (ALFA) and maximum residual in the solu- tion for the system of the two Helmholz equations (SYS), Helmholtz equation for the mean field equa- tion (M), solution of Z from ψ (Z) and solution of ψ from Z (PSI). To be defined by a DATA statement.

Q,FOCEAN,FCONT	real	The Helmholtz term for the mean field equation, friction coefficients over ocean and land. To be defined by a DATA statement.
WGT1,WGT2,WGT3	real	Weights for mixing near the boundary of boundary fields with forecasted fields. To be defined by a DATA statement.
ADIFF	real	Diffusion coefficient for humidity; it is defined by a DATA statement

Parameters only specified in Data statements.

STAB1	real	Static stability of layer 1
STAB2	real	Static stability of layer 2
PNIVS*	real	Pressure level P_0
PNIVM	real	Pressure level P_m
PNIV1	real	Pressure level P_1
IVAR	integer	Indicates if variable lateral boundaries are used; IVAR=0, constant boundary values; IVAR=1, variable boundary values.
KIND	integer	Indicates if channel is used (cyclic boundary conditions); KIND=1, channel; KIND=0, no channel (polar stereographic projection).

Subroutine STEP3P (F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,MY,F,MARK,
IT,IVAR,M,N)

The subroutine performs a computation for a time interval (6 hours) by the 3P-model including humidity and precipitation prediction. (See Flow Diagram B-2.)

<u>Subroutine Parameters</u>	<u>Description</u>
F1-F10	Working fields.
MY	μ -parameter field.
F	Coriolis parameter field.
MARK	Field indicator field.
IT	Array to store number of iterations for every timestep.
IVAR	Indicator for constant or variable boundary conditions; IVAR = 0 constant boundary; IVAR = 1 variable boundary.
M,N	Field vector.

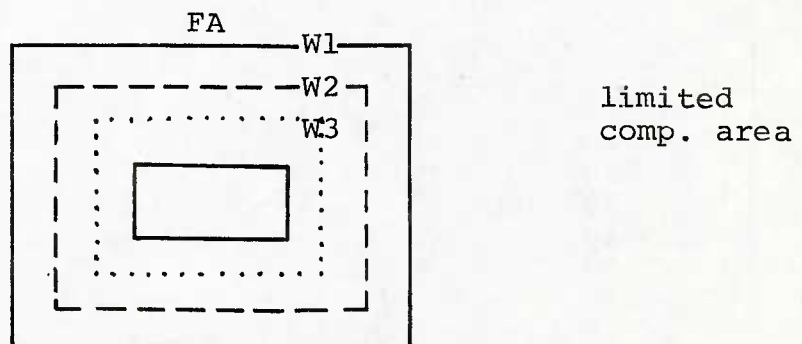
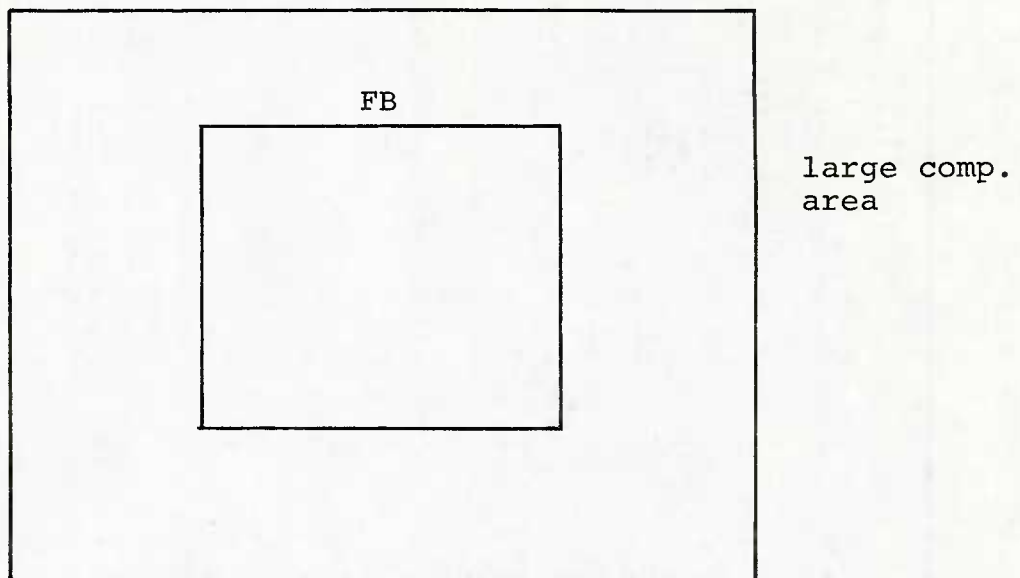
Parameters specified only in DATA statements

RGAS	$R = 287$
EE	$E = 0.622$
HL	$L = 2.5 \cdot 10^6$
CP	$C_p = 1004$
TO	$T_o = 273$
EO	$E_o = 0.611$
DEL1	δ_1
DEL2	δ_2
TOL	Tolerance

For explanation see chapter 8.4.

Subroutine MARK

The subroutine mixes field FA (corresponding to limited area) with FB (corresponding to a field taken from a large area).



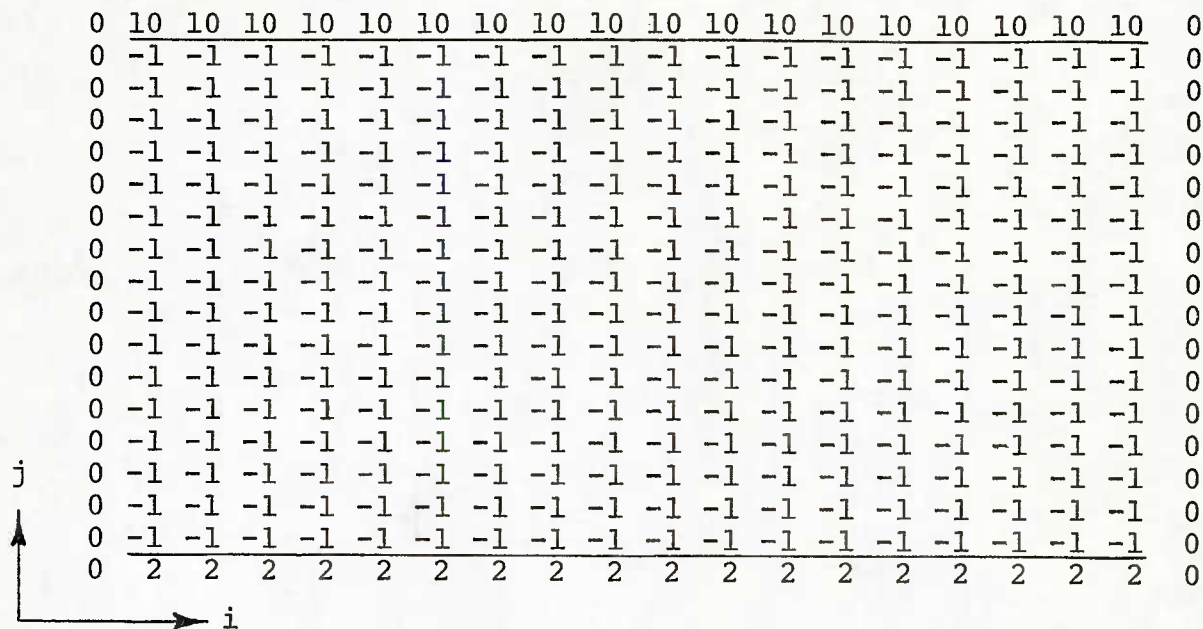
W_1, W_2, W_3
MN

Weight factors
Field vectors

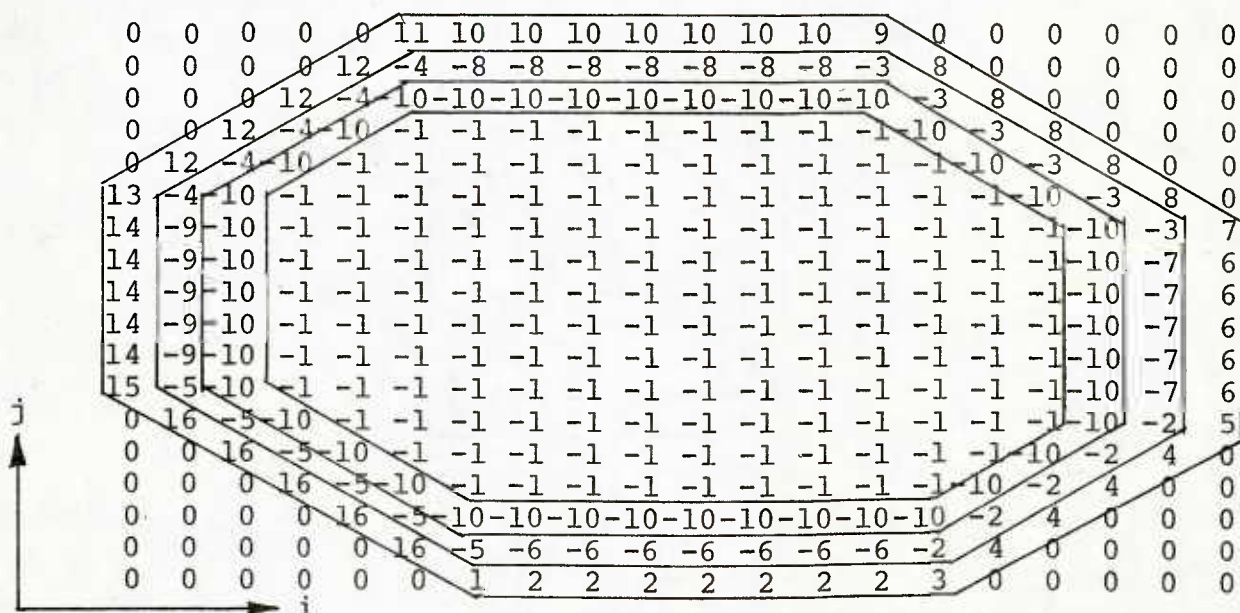
MARK-field

The MARK-field generated by the subroutine MARKF can be described by the following examples.

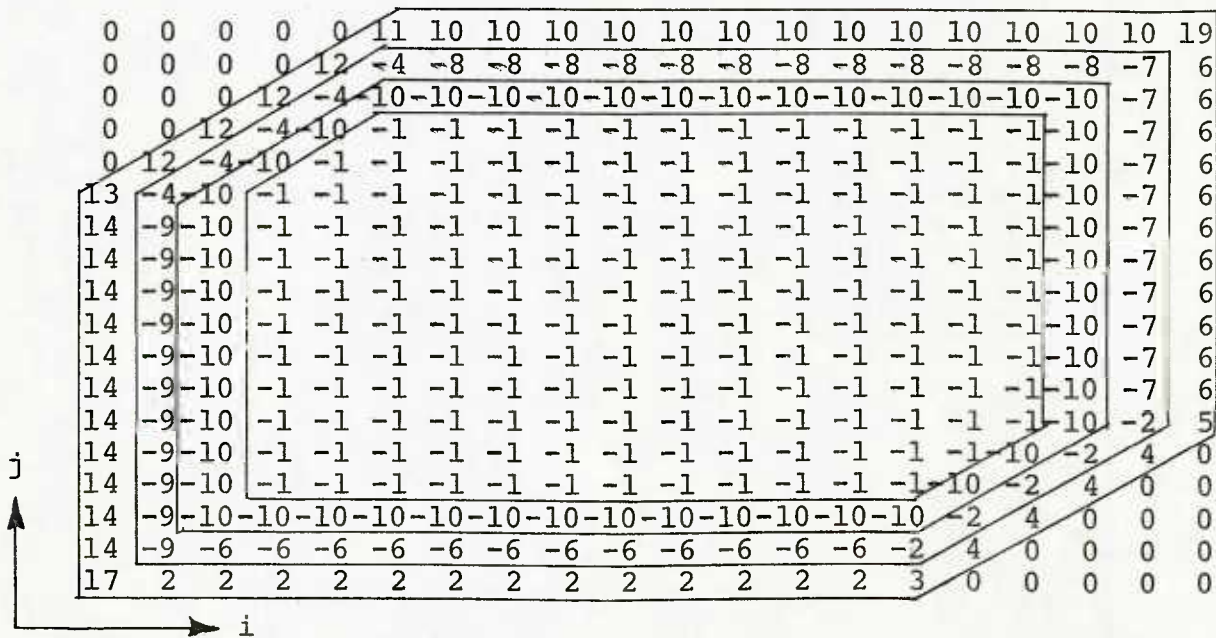
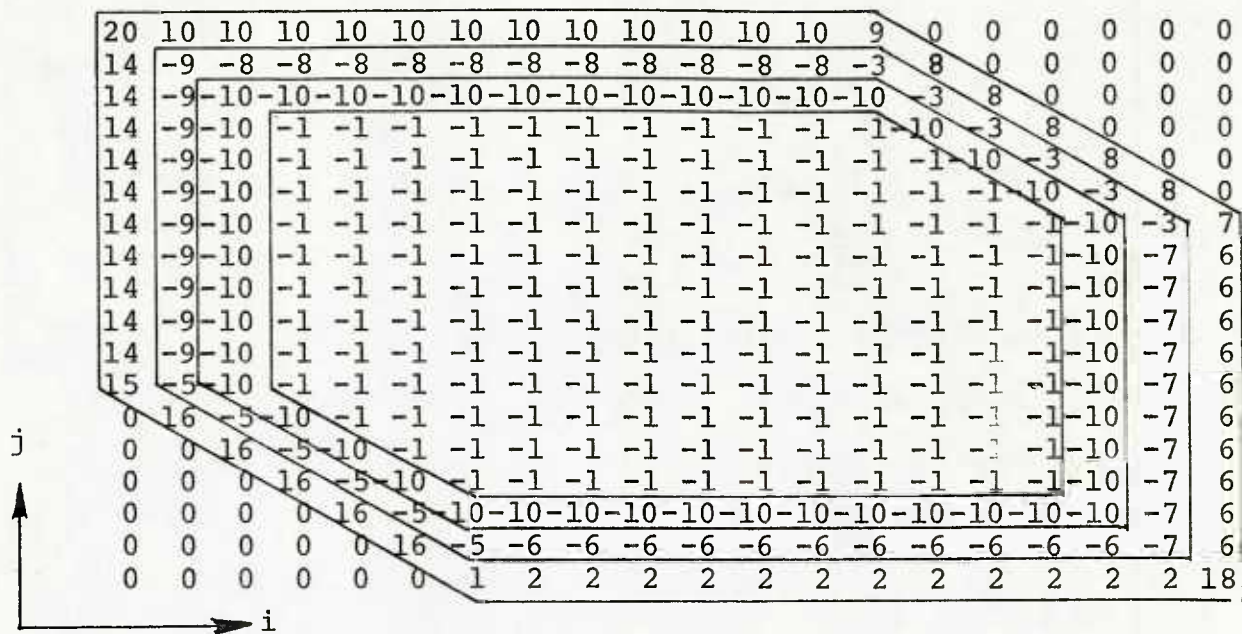
1. Channel field.



2. Ordinary octagonal field.



3. Fields with 2 rectangular points.



Subroutine RANWT (A,ALFA)

This subroutine writes field A (in core) to field ALFA (in secondary storage).

Subroutine RANRD (A,ALFA)

This subroutine reads field ALFA (in secondary storage) into field A (in core).

Subroutine MIXF(FA,FB,MARK,W1,W2,W3,M,N)

This subroutine mixes field FA (corresponding to limited area) with FB (corresponding to a field taken from a large area).

Subroutine HELMSYS (Z1,Z2,FORC1,FORC2,F,MY,FMY,A1,A2,B1,B2,
ALFA,RESIDUE,IT,M,N)

This subroutine solves the following system of Helmholtz equations by relaxation:

$$\nabla^2(Z1) - A1 \frac{f^2}{\mu} (Z1) + A2 \frac{f^2}{\mu} (Z2) = \text{FORC1}$$

$$\nabla^2(Z2) - B2 \frac{f^2}{\mu} (Z2) + B1 \frac{f^2}{\mu} (Z1) = \text{FORC2}$$

Subroutine
Parameter

Description

Z1,Z2

Fields to be solved by relaxation.
First guesses in Z1,Z2 before
using subroutine.

FORC1,FORC2

Forcing functions.

F

Coriolis parameter field.

MY

μ -parameter field.

FMY

Working field.

A1,A2,A3,A4

Physical parameters (constants).

ALPHA

α -overrelaxation coefficient.

RESIDUE

Residual, R, in the solution
of the system.

IT

Number of iterations.

M,N

Field vectors.

Subroutine MAP3P(I1,I2,I3,I4,I5,I6,I7,I8,I9,I10,I11,F1,F2,
F3,F4,F5,F6,F7,F8,F9,F10,F,MARK,M,N,IDAY,ITIME,NTIME,PNIVS,
PNIVM,PNIV1,ZB1,ZB2,ZB3)

Printing on line-printer of forecast fields in zebra patterns. Heights are computed from stream functions.

<u>Subroutine Parameters</u>	<u>Description</u>
I1-I11	Indicators. I = 0: no printing I = 1: printing. The numbers refer to the following fields: I1: Surface pressure. I2: Height for level p_m . I3: Height for level p_l . I4: Thickness ($p_m - p_s$). I5: Lower vertical velocity $\overline{\omega}_1$. I6: Precipitable water. I7: Accumulated precipitation. I8: Relative humidity. I9: Stream function for level p_s . I10: Stream function for level p_m . I11: Stream function for level p_l .
F1-F10	Working fields. (See flow diagram B-2)
F	Coriolis parameter field.
MARK	Indicator field (see MARKF) .
M,N	Field dimension.
IDAY	Year, month and day as one integer.
ITIME,NTIME	Initial and forecast time.
PNIVS,PNIVM,PNIV1	Pressure levels in the model.
ZB1,ZB2,ZB3	Working arrays for boundary strings.

Flow Diagram B-2: MAP3P

Variables to and
from secondary
storage

{PSIM1,PSI11
{PSI21

{WS,DIV1,DIV2
{HUM1,PREC

{ZM2,Z12,Z22

F1	F2	F3	F4	F5	F6	F7	F8	F9	F10
$\psi(P_S)$.	$\psi(P_m)$.	$\psi(P_L)$.	ω_S x $\frac{\omega_L}{\omega_L}$	D_L x	D_2 x	q x	ψ_m x Σr x	ψ_L x	ψ_2 x
x	x	x	.	$Z(P_S)$ x	$Z(P_m)$ x	$Z(P_L)$.	Z_m^T .	Z_L^T .	Z_2^T .
		p	27_L .						

MAP

MAP

STREAMF

MAP

Subroutine STREAMF (Z,PSI,R,F,MARK, ZB,PSIB,FB,SB,GAMMA,IND,
IT,M,N)

The subroutine computes the linearized balance equation.

<u>Subroutine Parameter</u>	<u>Description</u>
Z	Z field.
PSI	PSI field.
R	Forcing function for the Poisson equation in the solution of $\nabla^2 \psi = F(Z)$ or $\nabla^2 Z = Z(\psi)$.
F	Coriolis parameter.
MARK	MARK-field.
ZB	String of boundary values for Z_k .
PSIB	String of boundary values for ψ_k .
FB	String of boundary values for f_k .
SB	String of boundary values for $(\Delta S)_k$.
GAMMA	γ .
IND	See subroutine comments.
IT	Number of iterations for solving the Poisson equation by relaxation.
M,N	Field vectors.

Subroutines called by STREAMF

BDRGRD(SB)	Computes SB.
BRDVAL(A,AB,M,IND)	Computes boundary values from field A and stores the boundary values in string AB or reverse: IND = 0, AB = A; IND = 1, A = AB.
M	Field parameter
IND	See subroutine comments

POIS is an entry point in HELM for the solution of a Helmholtz equation by Liebmann relaxation.

HELM(Z,FORC,Q,ALFA,RESIDUE,IT,M,N)

<u>Subroutine Parameter</u>	<u>Description</u>
Z	Z-field.
FORC	Forcing function.
Q	Helmholtz coefficient.
ALFA	Overrelaxation coefficient.
RESIDUE	Residual.
M,N	Field vectors.
IT	Counts the number of iterations required to obtain a solution.

Subroutine SATUR(TM, QSAT)

This subroutine computes QSAT, integrated mixing ratio at saturation as a function of the mean temperature between 500 and 1000 mb. TM values of QSAT are given for each whole degree between -50°C to +20°C in the subroutine. Linear interpolation between whole degrees is employed and the result QSAT is given in $\text{TON/m}^2/\text{cb}$.

Subroutine MAP(Q,M,N,DS,QZ,QD,SCALE)

Prints a field in "zebra pattern" on line-printer.

Subroutine Parameters

Description

Q

Field to be printed.

M,N

Field vectors.

DS

Grid distance (meters)

QZ

Isoline corresponding to 000, the
line towards 999 on printer
(indicated by heavy line below).

QD

Resolution indicator (see below).

000000000 } QD
000000000 }
 } QD

111111111
111111111

222222222
222222222

SCALE

Map scale factor.
(unit 10^6 m)

Subroutine JACOB(A,B,C,M,N)

This subroutine computes a Jacobian operator (of the form)

$$C = J(A,B) = (A_{i+1} - A_{i-1})_j (B_{j+1} - B_{j-1})_i \\ - (A_{j+1} - A_{j-1})_i (B_{i+1} - B_{i-1})_j .$$

<u>Subroutine parameters</u>	<u>Description</u>
A,B,C	Fields according to formula.
M,N	Field vectors .

Subroutine ABSVOR(PSI,VOR,F,MY,MARK,M,N)

This subroutine computes the relative or the absolute vorticity.

- a. $\eta = \mu \nabla^2 \psi + f$ ABSVOR
b. $\zeta = \mu \nabla^2 \psi$ RELVOR (via entry point)

<u>Subroutine Parameters</u>	<u>Description</u>
PSI	ψ field .
VOR	η or ζ field .
F	Coriolis parameter field .
MY	μ parameter field .
MARK	Field indicator array.
M,N	Field vectors .

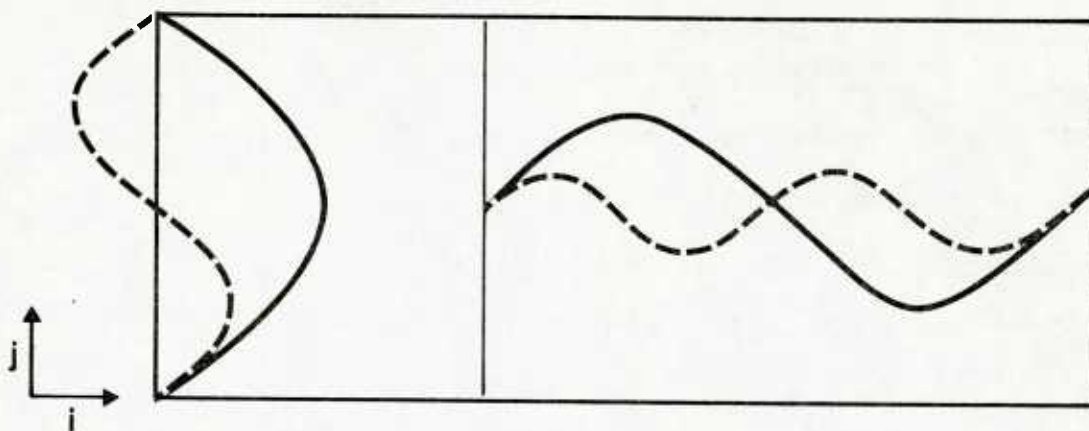
Subroutine GENCH(PSI,M,N,PSIO,U,NU,NWAVE,NX,NY,PSIC,LAMC,
PSIS,LAMS,WX)

<u>Subroutine Parameters</u>	<u>Description</u>
PSI	Result field.
M,N	Field vectors.
PSIO	Constant value (ψ_0).
U	Zonal wind speed.
NU	Resolution of zonal wind speed.
NWAVE	Number of waves.
NX	Wave numbers as a function of channel length in x-direction (nx).
NY	Wave numbers as a function of channel width in y-direction (ny).
PSIC	Amplitudes of the cosine function (ψ_c).
LAMC	Phase differences for the cosine functions (ψ_c) in whole degrees.
PSIS	Amplitudes of the sine functions (ψ_s) in whole degrees.
LAMS	Phase differences for the sine functions (λ_s) in whole degrees.
WX	Adjustment of the wave near the rigid boundaries: Boundary, WX = 0; +1 row from the boundary, WX = 0.33; +2 row from the boundary, WX = 0.64; +3 and more, WX = 1.

Fields are generated according to the formula:

$$\psi = \psi_0 - Uy + \sum_{v=1}^N [(\psi_c)_v \sin\{(nx)_v(x) - \frac{\pi}{180}(\lambda_c)_v\} \cos(ny)_v + (\psi_s)_v \sin\{(nx)_v(x) - \frac{\pi}{180}(\lambda_s)_v\} \sin(ny)_v]$$

Example Prediction Area



- Nx = 1 1 wave in x-direction over the area (solid line)
- Nx = 2 2 waves in x-direction over the area (dashed line)
- Nx = N N waves in x-direction over the area
- Ny = 1 1 half wave in y-direction (solid line)
- Ny = 2 2 half waves in y-direction (dashed line)
- Ny = N N half waves in y-direction

The zonal wind speed is specified with an arbitrary resolution across the channel. All parameters are given by a DATA statement.

Subroutine STEPEXT(F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,MY,F,MARK,
M,N,I1,I2,I3,I4)

The subroutine computes the forcing functions

$$\sum_{i=1}^4 F_{mi}, \quad \sum_{i=1}^4 F_{li}, \quad \text{and} \quad \sum_{i=1}^4 F_{2i}$$

and stores the result in the fields HM3, H13 and H23
 respectively on secondary memories (see Flow Diagram B-3).

The program is a subroutine to STEP3P.

<u>Subroutine Parameters</u>	<u>Description</u>
F1-F10	Working field.
MY	μ -parameter field.
F	Coriolis parameter field.
MARK	Field indicator array.
I1,I2,I3,I4	Computational parameters.
	$I1 = 0 \quad W_{\chi} \cdot \nabla \eta = 0$
	$I1 = 1 \quad W_{\chi} \cdot \nabla \eta \neq 0.$
	$I2 = 0 \quad \zeta \cdot \nabla W = 0.$
	$I2 = 1 \quad \zeta \cdot \nabla W \neq 0.$
	$I3 = 0 \quad \omega \frac{\partial \zeta}{\partial p} = 0.$
	$I3 = 1 \quad \omega \frac{\partial \zeta}{\partial p} \neq 0.$
	$I4 = 0 \quad k \cdot (\frac{\partial V}{\partial p} \times \nabla \omega) = 0.$
	$I4 = 1 \quad k \cdot (\frac{\partial V}{\partial p} \times \nabla \omega) \neq 0.$
M,N	Field vectors.

Flow Diagram B-3: STEPEXT

F1	F2	F3	F4	F5	F6	F7	F8	F9	F10	To second storage	From second storage	Comments
ψ_m^T	ψ_1^T	ψ_2^T	W_B	Heat	-	-	-	J_4		J4=F8	F5=DIV1 F6=DIV2	
				D_1^{-1}	D_2^{-1}							
						γ_{10}	γ_{11}	$\psi_1 \psi_{10}$	$\psi_2 \psi_{11}$			
								F14	F24	H13=F9 H23=F10		Twisting term for thickness fields to second storage
						γ_7	γ_8	$\psi_1 \psi_7$	$\psi_2 \psi_8$			
						γ_9	$\psi_9 \psi_9$					
							F_m^4			H43=F8		Twisting term for mean field to second storage.
						ζ_m	ζ_1	ζ_2				
F_{m2}	F_{12}	F_{22}										
$F_{m2}^* F_{m3}$	$F_{12}^* F_{13}$	$F_{22}^* F_{23}$								H42=F1 H12=F2 H22=F3		$J_5 = \psi \cdot V$ to secondary storage
$\psi^2 \chi_m$	$\psi^2 \chi_1$	$\psi^2 \chi_2$	$(\chi_m)_g$	$(\chi_1)_g$	$(\chi_2)_g$					F4=V1 F5=V2 F6=V2		Compute forcing function for velocity potential in order to get velocity
$\eta_m - 2\zeta$	$\eta_m - 2\zeta_2$		χ_m	χ_1	χ_2					V4=F4 V1=F5 V2=F6		
									$\psi_1 \psi_{13}$			$\gamma_{13} = \eta_m - 2\zeta_1$
									F_{11}	H11=F10		$\psi_1 \eta$ to secondary storage
									$\psi_2 \psi_{14}$			$\gamma_{14} = \eta_m - 2\zeta_2$
δ_7	δ_8	δ_9							F_{21}	H21=F10		$\psi_1 \eta$ to secondary storage
F_{m1}						$\psi_m \psi_{d7}$	$\psi_1 \psi_8$	$\psi_2 \psi_9$				
$F_{m2}^* F_{m3}$	F_{m4}	F_{11}	$F_{12}^* F_{13}$	F_{14}	F_{21}	$F_{22}^* F_{23}$	F_{24}					
IF_{m1}		IF_{11}			IF_{21}					H43=F1 H13=F4 H23=F7		Sum of higher order terms to H43, H13 and H23
0			0		0							
ψ_m^T	ψ_1^T	ψ_2^T	W_B	Heat				J_4				Restore initial fields

Subroutine VELPOT(KSI, FORC, M, N, RESIDUE, ALFA)

This subroutine computes the velocity potential from a known divergence field

$$\nabla^2 \chi = D.$$

In finite-difference form

$$\nabla^2 \chi = \frac{D}{\mu}.$$

The solution is performed by Liebmann relaxation with an overrelaxation coefficient ALFA equal approximately to 1.4, but its size depends on the area and mesh width. The residual RESIDUE must also be given ($0.5 \cdot 10^{+6}$ recommended value).

<u>Subroutine Parameters</u>	<u>Description</u>
KSI	2D array for the χ field.
FORC	2D array for the forcing function.
ALFA	Overrelaxation coefficient.
RESIDUE	Residual.
M, N	Field vectors.

Remark: A first guess must be put in χ field before the execution.

Subroutine GRADPR(A,B,C,MARK,M,N)

This subroutine computes a finite difference operation of the form $\nabla A \cdot \nabla B$.

$$\begin{aligned} (\nabla A \nabla B)_{ij} = & (A_{i+1} - A_i) (B_{i+1} - B_i)_j + (A_i - A_{i-1}) (B_i - B_{i-1})_j \\ & + (A_{j+1} - A_j) (B_{j+1} - B_j)_i + (A_j - A_{j-1}) (B_j - B_{j-1})_i \end{aligned}$$

<u>Subroutine Parameters</u>	<u>Description</u>
A,B	Fields for operational vector.
C	Result field.
MARK	Field indicator array.
M,N	Field vectors.

APPENDIX C
PROGRAM LISTINGS

Three programs for the three-parameter model in Fortran IV are presented: PROG3P, STEP3P, and STEPEXT. The other subroutines may be obtained from ENVPREDRSCHFAC by request.

PROGRAM PROG3P

```

PROGRAM PROG3P(OUTPUT,TAPE6=OUTPUT,DATA,INPUT=DATA)
* BAROCLINIC BALANCED INTEGRATED 3-PARAMETER MODEL INCLUDING HUMIDITY
* AND PRECIPITATION. BOUNDARY VALUES CAN BE VARIABLE, CONSTANT OR OF
* CHANNEL TYPE.
  DIMENSION F1(57,57),F2(57,57),F3(57,57),F4(57,57),F5(57,57),
1          F6(57,57),F7(57,57),F8(57,57),F9(57,57),F10(57,57),
2          F(57,57),MY(57,57),MARK(57,57)
  DIMENSION UPS(10),UPM(10),UP1(10),WX(15),NX(20),NY(20),PSIC(20),
X          PSIS(20),LAMC(20),LAMS(20)
  DIMENSION ZB(250),PSIB(250),FB(250),SB(250),IT(2,25)
  DIMENSION MAPHOUR(10),LETACC(10),NSMUTT(10),NELLIPT(10)
  DIMENSION MSMUTT(10)
  DIMENSION MELLIPT(10)

COMMON/FPAR/IC(8),JC(8),XPOL,YPOL,R,RE,DS,JMIN(100),JMAX(100),
X          M,N,KIND
COMMON/FORM1/IC1(8),JC1(8),JMIN1(100),JMAX1(100),XP1,YP1,D1
COMMON/DRM/MN,NDIM,NFLD
COMMON/KANAL/FIM
COMMON F
COMMON/ECS/ PSIM1,PSI11,PSI21,PSIM2,PSI12,PSI22,HUM1,HUM2,DIV1,
2DIV2,WS,HEAT,J789,J12,J56,J3,PS,TS,PREC,STRM,STR1,STR2,ZM1,Z11,Z21
3,ZM2,Z12,Z22,H13,H23,HM3,HM2,H12,H22,H11,H21,J4,VM,V1,V2
COMMON/COEFF/A1,A2,A3,B1,B2,B3,C1,C2,C3,C4,C5,C6,C7,C8,D,DELP,EM,
X          E1,E2,H1,H2,H3,H4,H5,H6,PMEAN,S1,S2,T1,T2,T3,T4,T5,
X          P0,PM,P1
COMMON/COEFF2/T6,T7,T8,T9,T10,T11,T12,T13,T14,
X          K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15,
X          K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28,
X          K29,K30,K31,K32,K33,K34,K35,K36,K37,K38
COMMON/RUNPAR/DELT,NTSTEP,ALFASYS,ALFAM,ALFAZ,ALFAPSI,RESSYS,RESM,
X          RESZ,RESPSI,Q,FOCEAN,FCONT,WGT1,WGT2,WGT3,ADIFF
REAL MY
REAL          K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15,
X          K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28,
X          K29,K30,K31,K32,K33,K34,K35,K36,K37,K38
DATA PSIM1,PSI11,PSI21,PSIM2,PSI12,PSI22,HUM1,HUM2,DIV1,
*DIV2,WS,HEAT,J789,J12,J56,J3,PS,TS,PREC,STRM,STR1,STR2,ZM1,Z11,Z21
*,ZM2,Z12,Z22,H13,H23,HM3,HM2,H12,H22,H11,H21,J4,VM,V1,V2
*/1,3250,6499,9748,12997,16246,19495,22744,25993,29242,32491,35740,
*38989,42238,45487,48736,51985,55234,58483,61732,64981,68230,
*71479,74728,77977,81226,84475,87724,90973,94222,97471,100720,
*103969,107218,110467,113716,116965,120214,123463,126712/
NAMELIST/FPARR/IC,JC,XPOL,YPOL,R,RE,DS,JMIN,JMAX,M,N,KIND
DATA(IC1(1),I=1,8)/1,1,57,57,57,57,1,1/
DATA(JC1(1),I=1,8)/1,1,1,1,57,57,57,57/
DATA(XP1,YP1,D1)/-21,,29,,95,25/
DATA FIM/ 50, /
DATA KIND/0/
DATA ICASE,IDAY,ITIME/1,660922,00/
DATA ALFASYS,ALFAM,ALFAZ,ALFAPSI/,85,1.45,1,4,1,8/
DATA RESSYS,RESM,RESZ,RESPSI/,5E5,,5,,5E5 /
DATA Q,FOCEAN,FCONT,WGT1,WGT2,WGT3,ADIFF/ 0,0 ,.62,1,,1,,5,,2,
*          0,0/
DATA STAB1,STAB2,PNIVS,PNIVM,PNIV1/,422222,,511111,100,,50,,30./
C NTSTEP IS NUMBER OF TIME STEPS IN 3 HOURS
DATA NTSTEP/4/
DATA NEND/48/
DATA IVAR/ 0 /

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PROG3P (Continued)

C 0 AND ADIFF HAVE TO BE DEFINED ***	PROG3P.60
C MAPHOUR GIVES THE HOURS(MULTIPLE OF 6) FOR MAP-PRINTING.	PROG3P.61
C LETACC GIVES THE HOURS(MULTIPLE OF 6) AT WHICH THE ACCUMULATED PRECIPITATION IS PRINTED.	PROG3P.62
C SHALL BE PUT =0 AGAIN.	PROG3P.63
C NSMUTT GIVES THE HOURS(MULTIPLE OF 6) FOR SMOOTHING.	PROG3P.64
C NELLIPT GIVES THE HOURS (MULTIPLE OF 6) FOR ELLIPTICITY TEST.	PROG3P.65
DATA MAPHOUR/0,12,24,36,48,-1,-1,-1,-1/	JUN12.6
DATA (LETACC(I),I=1,10)/0,12,24,36,48,-1,-1,-1,-1/	JUN12.7
DATA(NSMUTT(I),I=1,10)/0,3,6,9,12,15,18,21,24,-1/	PROG3P.68
DATA (MSMUTT(I),I=1,10)/27,30,33,36,39,42,45,48,-1,-1/	JUN12.8
DATA(NELLIPT(I),I=1,10)/0,3,6,9,12,15,18,21,24,-1/	PROG3P.69
DATA (NELLIPT(I),I=1,10)/27,30,33,36,39,42,45,48,-1,-1/	JUN12.9
3=9,806	PROG3P.70
F0=1.03E-4	PROG3P.71
C TIME STEP IN SECONDS	APR14.14
DELT=1.*3.6E3/NTSTEP	JUN12.10
LABEL = 10H PUT1	PROG3P.73
BTIME = SECOND(FAKE)	PROG3P.74
WRITE(6,8000) BTIME	PROG3P.75
8000 FORMAT(1H0, *BTIME=*, F10.4)	PROG3P.76
CALL PUT1	PROG3P.77
DTIME = SECOND(FAKE) - BTIME	PROG3P.78
WRITE(6,8005) LABEL, DTIME	PROG3P.79
8005 FORMAT(1H0, *TIME TO EXECUTE*, A10, F10.4)	PROG3P.80
LABEL = 10H COEFF3P	PROG3P.81
BTIME = SECOND(FAKE)	PROG3P.82
WRITE(6,8000) BTIME	PROG3P.83
CALL COEFF3P(STAB1,STAR2,PNIVS,PNIVM,PNIV1)	PROG3P.84
DTIME = SECOND(FAKE) - BTIME	PROG3P.85
WRITE(6,8005) LABEL, DTIME	PROG3P.86
MN = M*N	PROG3P.87
NDIM = 20000	PROG3P.88
LABEL = 10H MARKF	PROG3P.89
BTIME = SECOND(FAKE)	PROG3P.90
WRITE(6,8000) BTIME	PROG3P.91
CALL MARKF(MARK,M,N)	PROG3P.92
DTIME = SECOND(FAKE) - BTIME	PROG3P.93
WRITE(6,8005) LABEL, DTIME	PROG3P.94
LABEL = 10H MYFF	PROG3P.95
BTIME = SECOND(FAKE)	PROG3P.96
WRITE(6,8000) BTIME	PROG3P.97
CALL MYFF(MY,M,N)	APR14.16
DTIME = SECOND(FAKE) - BTIME	PROG3P.98
WRITE(6,8005) LABEL, DTIME	PROG3P.100
DO 9 I=1,MN	PROG3P.101
9 F1(I)=0.0	PROG3P.102
LABEL = 10H RANWT	PROG3P.103
BTIME = SECOND(FAKE)	PROG3P.104
WRITE(6,8000) BTIME	PROG3P.105
CALL RANWT(VM,F1)	PROG3P.106
CALL RANWT(V1,F1)	PROG3P.107
CALL RANWT(V2,F1)	PROG3P.108
DTIME = SECOND(FAKE) - BTIME	PROG3P.109
WRITE(6,8005) LABEL, DTIME	PROG3P.110
VTIME = 0	PROG3P.111
IF(KIND,NE,0) GO TO 45	PROG3P.112
C INITIAL FIELDS. NO CHANNEL.	PROG3P.113
C THE SUBROUTINE INITF HAS TO BE WRITTEN ***	PROG3P.114
C INITIAL Z-FIELDS AT THE LEVELS PS,PM AND P1 TO F1,F2 AND F3. THE HUMIDITY F4	PROG3P.115
C FIELD Q TO F4. STANDARD SURFACE PRESSURE PS TO F5 AND SEA LEVEL TEMPERATURE F6	PROG3P.116

PROG3P (Continued)

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C TS TO F6, DEFINITION Q=0.0334892*M(850)+0.0363689*M(700)+0.0290951*M(600)+0.0145476*M(300) WHERE M(P) ARE MIXING RATIOS IN MTS=UNITS, PS=101.325
C 0.0145476*M(300) WHERE M(P) ARE MIXING RATIOS IN MTS=UNITS, PS=101.325
C AT SEA SURFACE.
    LABEL = 10H INITF
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
    CALL INITF(F1,F2,F3,F4,F5,F6,F7(1,1),F7(1,5),F7(1,9),F7(1,13),
    *F7(1,17),F7(1,21),F7(1,25),F8,M,N)
    SCALE=20.
    DTIME = SECOND(FAKE) - BTIME
    WRITE(6,8005) LABEL, DTIME
    LABEL = 10H MAP 6X
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
C SOLUTION OF STREAMFUNCTIONS
    LABEL = 10H STREAMF 1
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
    LABEL = 10H STREAMF
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
15 CALL STREAMF(F1,F7,F10,MARK,ZB,PSIB,FB,SB,GAMMA,0,IT1,M,N)
    DTIME = SECOND(FAKE) - BTIME
    WRITE(6,8005) LABEL, DTIME
    LABEL = 10H STREAMF 2
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
    CALL STREAMF(F2,F8,F10,MARK,ZB,PSIB,FB,SB,GAMMA,0,IT2,M,N)
    DTIME = SECOND(FAKE) - BTIME
    WRITE(6,8005) LABEL, DTIME
    LABEL = 10H STREAMF 3
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
    CALL STREAMF(F3,F9,F10,MARK,ZB,PSIB,FB,SB,GAMMA,0,IT3,M,N)
    DTIME = SECOND(FAKE) - BTIME
    WRITE(6,8005) LABEL, DTIME
    DO 20 I=1,MN
    F5(I)=F5(I)*0.1
    Z=F2(I)
    F2(I)=.5*(Z+F1(I))
    F3(I)=.5*(F3(I)-Z)
    F1(I)=Z
20 CONTINUE
C ADAPTION OF HUMIDITY FIELD
    DO 21 I=1,MN
    TZ = G*(H3*F2(I)+H4*F3(I))/F0
    TPSI = .5*(H3*(F8(I)-F7(I))+H4*(F9(I)-F8(I)))
    CALL SATUR(TZ,QZ)
    CALL SATUR(TPSI,QPSI)
C CHANGE IN F4 SINCE F4 IS RELATIVE HUMIDITY
    F4(I)=F4(I)*QPSI
C
    F4(I)=F4(I)*QPSI/QZ
    Z = F4(I)-.8*QPSI
    IF(Z.GT.0.0) F4(I)=.8*QPSI
21 CONTINUE
C STORE Z-FIELDS FOR BOUNDARY MIXING
    LABEL = 10H RANWT 3X
    BTIME = SECOND(FAKE)
    WRITE(6,8000) BTIME
    CALL RANWT(ZM2,F1)
    CALL RANWT(Z12,F2)

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PROG3P.117
 PROG3P.118
 PROG3P.119
 PROG3P.120
 PROG3P.121
 PROG3P.122
 PROG3P.123
 APR14,17
 PROG3P.125
 PROG3P.126
 PROG3P.127
 PROG3P.128
 PROG3P.129
 PROG3P.130
 PROG3P.139
 PROG3P.140
 PROG3P.141
 PROG3P.142
 PROG3P.143
 PROG3P.144
 PROG3P.145
 APR14,18
 PROG3P.147
 PROG3P.148
 PROG3P.149
 PROG3P.150
 PROG3P.151
 APR14,19
 PROG3P.153
 PROG3P.154
 PROG3P.155
 PROG3P.156
 PROG3P.157
 APR14,20
 PROG3P.159
 PROG3P.160
 PROG3P.161
 PROG3P.162
 PROG3P.163
 PROG3P.164
 PROG3P.165
 PROG3P.166
 PROG3P.167
 PROG3P.168
 PROG3P.169
 PROG3P.170
 PROG3P.171
 PROG3P.172
 PROG3P.173
 PROG3P.174
 PROG3P.175
 PROG3P.176
 PROG3P.177
 PROG3P.178
 PROG3P.179
 PROG3P.180
 PROG3P.181
 PROG3P.182
 PROG3P.183
 PROG3P.184
 PROG3P.185

PROG3P (Continued)

CALL RANWT(Z22,F3)	PROG3P.186
DTIME = SECOND(FAKE) = BTIME	PROG3P.187
WRITE(6,8005) LABEL, DTIME	PROG3P.188
C TEST FOR ELLIPTICITY OF PSI-FIELDS	PROG3P.189
LABEL = 10H ELLIPT 3X	PROG3P.190
BTIME = SECOND(FAKE)	PROG3P.191
WRITE(6,8000) BTIME	PROG3P.192
CALL ELLIPT(F7,MY,M,N)	APR14,21
CALL ELLIPT(F8,MY,M,N)	APR14,22
CALL ELLIPT(F9,MY,M,N)	APR14,23
DTIME = SECOND(FAKE) = BTIME	PROG3P.196
WRITE(6,8005) LABEL, DTIME	PROG3P.197
C COMPUTATION OF PSIM,PSI1 AND PSI2	PROG3P.198
DO 25 I=1,MN	PROG3P.199
Z=F8(I)	PROG3P.200
F8(I)=.5*(Z-F7(I))	PROG3P.201
F9(I)=.5*(F9(I)-Z)	PROG3P.202
F7(I)=Z	PROG3P.203
25 CONTINUE	PROG3P.204
IF(IVAR.EQ.0) GO TO 31	PROG3P.205
C STORE FIELDS FOR BOUNDARY MIXING	PROG3P.206
DO 30 I=1,MN	PROG3P.207
Z=F1(I)	PROG3P.208
F1(I)=F7(I)	PROG3P.209
F7(I)=F7(I)-G*Z/F0	PROG3P.210
Z=F2(I)	PROG3P.211
F2(I)=F8(I)	PROG3P.212
F8(I)=F8(I)-G*Z/F0	PROG3P.213
Z=F3(I)	PROG3P.214
F3(I)=F9(I)	PROG3P.215
F9(I)=F9(I)-G*Z/F0	PROG3P.216
30 CONTINUE	PROG3P.217
GO TO 36	PROG3P.218
31 DO 35 I=1,MN	PROG3P.219
F1(I)=F7(I)	PROG3P.220
F2(I)=F8(I)	PROG3P.221
F3(I)=F9(I)	PROG3P.222
35 CONTINUE	PROG3P.223
LABEL = 10H RANWT 3X	PROG3P.224
BTIME = SECOND(FAKE)	PROG3P.225
WRITE(6,8000) BTIME	PROG3P.226
36 CALL RANWT(STR1,F7)	PROG3P.227
CALL RANWT(STR1,F8)	PROG3P.228
CALL RANWT(STR2,F9)	PROG3P.229
DTIME = SECOND(FAKE) = BTIME	PROG3P.230
WRITE(6,8005) LABEL, DTIME	PROG3P.231
C PRINT NUMBER OF ITERATIONS IN THE SOLUTION OF PSI	PROG3P.232
41 WRITE(6,202) IT1,IT2,IT3	PROG3P.233
202 FORMAT(///1X,33HNUMBER OF ITERATIONS IN STREAMF ,3(14))	PROG3P.234
GO TO 53	PROG3P.235
C GENERATION OF INITIAL FIELDS FOR THE CHANNEL VERSION, PSI-FIELDS ARE	PROG3P.236
C GENERATED IN THE SUBROUTINE GENCH FROM PARAMETERS GIVEN IN DATA-STATEM	PROG3P.237
C HUMIDITY-FIELD CORRESPONDS TO 50 PERCENT RELATIVE HUMIDITY. PS IS 100	PROG3P.238
C AND TS 273 DEGREES EVERYWHERE.	PROG3P.239
45 CALL GENCH(F7,M,N,PSIPS,UPS,NU,NWAVE,NX,NY,PSIC,LAMC,PSIS,LAMS,WX)	PROG3P.240
WRITE(6,FPARR)	PROG3P.241
CALL GENCH(F8,M,N,PSIPM,UPM,NU,NWAVE,NX,NY,PSIC,LAMC,PSIS,LAMS,WX)	PROG3P.242
CALL GENCH(F9,M,N,PSIP1,UP1,NU,NWAVE,NX,NY,PSIC,LAMC,PSIS,LAMS,WX)	PROG3P.243
C COMPUTATION OF PSIM,PSI1,PSI2,HUMIDITY,PS AND TS	PROG3P.244
46 DO 50 I=1,MN	PROG3P.245

PROG3P (Continued)

F1(I)=F8(I)	PROG3P,246
F2(I)=.5*(F8(I)+F7(I))	PROG3P,247
F3(I)=.5*(F9(I)+F8(I))	PROG3P,248
F5(I)= 100.	PROG3P,249
F6(I)= 273.	PROG3P,250
TPSI= H3*F2(I)+H4*F3(I)	PROG3P,251
CALL SATUR(TPSI,QPSI)	PROG3P,252
F4(I)=.5*QPSI	PROG3P,253
50 CONTINUE	PROG3P,254
C STORE PSI-FIELDS IN STRM,STR1 AND STR2	PROG3P,255
CALL RANWT(STRM,F1)	PROG3P,256
CALL RANWT(STR1,F2)	PROG3P,257
CALL RANWT(STR2,F3)	PROG3P,258
C SMOOTHING OF INITIAL FIELDS	PROG3P,259
53 IF(NSMUTT(1).NE.0) GO TO 55	PROG3P,260
LABEL = 10H ASMUT 6X	PROG3P,261
BTIME = SECOND(FAKE)	PROG3P,262
WRITE(6,8000) BTIME	PROG3P,263
CALL ASMUT(F1,F10,M,N,.5)	PROG3P,264
CALL ASMUT(F1,F10,M,N,-.5)	PROG3P,265
CALL ASMUT(F2,F10,M,N,.5)	PROG3P,266
CALL ASMUT(F2,F10,M,N,-.5)	PROG3P,267
CALL ASMUT(F3,F10,M,N,.5)	PROG3P,268
CALL ASMUT(F3,F10,M,N,-.5)	PROG3P,269
DTIME = SECOND(FAKE) - BTIME	PROG3P,270
WRITE(6,8005) LABEL, DTIME	PROG3P,271
C LOADING OF INITIAL FIELDS, ZERO TO ACCUMULATED PRECIPITATION	PROG3P,272
55 DO 56 I=1,MN	PROG3P,273
F10(I)=0.0	PROG3P,274
56 CONTINUE	PROG3P,275
LABEL = 10H RANWT 10X	PROG3P,276
BTIME = SECOND(FAKE)	PROG3P,277
WRITE(6,8000) BTIME	PROG3P,278
CALL RANWT(PSIM1,F1)	PROG3P,279
CALL RANWT(PSI11,F2)	PROG3P,280
CALL RANWT(PSI21,F3)	PROG3P,281
CALL RANWT(HUM1, F4)	PROG3P,282
CALL RANWT(PS, F5)	PROG3P,283
CALL RANWT(TS, F6)	PROG3P,284
CALL RANWT(PREC, F10)	PROG3P,285
CALL RANWT(WS, F10)	PROG3P,286
CALL RANWT(DIV1, F10)	PROG3P,287
CALL RANWT(DIV2, F10)	PROG3P,288
DTIME = SECOND(FAKE) - BTIME	PROG3P,289
WRITE(6,8005) LABEL, DTIME	PROG3P,290
IF(MAPHOUR(1).NE.0) GO TO 60	PROG3P,291
C MAPPRINTING OF FIRST TIMESTEP	PROG3P,292
IF(KIND.NE.0) GO TO 58	PROG3P,293
LABEL = 10H MAP3P	PROG3P,294
BTIME = SECOND(FAKE)	PROG3P,295
WRITE(6,8000) BTIME	PROG3P,296
CALL MAP3P(1,1,1,1,0,2,0,2,0,0,0,F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,	APR14,24
X MARK,M,N,IDAY,ITIME,NTIME,PNIVS,PNIVM,PNIV1,ZB,PSIR,FB)	PROG3P,298
DTIME = SECOND(FAKE) - BTIME	PROG3P,299
WRITE(6,8005) LABEL, DTIME	PROG3P,300
GO TO 60	PROG3P,301
58 CALL MAP3P(0,0,0,0,0,0,0,0,1,1,1,F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,	APR14,25
X MARK,M,N,IDAY,ITIME,NTIME,PNIVS,PNIVM,PNIV1,ZB,ZB,ZB)	PROG3P,303
LABEL = 10H FORECASTS	PROG3P,304
BTIME = SECOND(FAKE)	PROG3P,305

PROG3P (Continued)

WRITE(6,8000) BTIME	PROG3P,306
C THREE HOURS FORECAST STARTS HERE	PROG3P,307
60 DO 61 I=1,MN	PROG3P,308
F1(I)= 0.0	PROG3P,309
61 CONTINUE	PROG3P,310
CALL RANWT(DIV1,F1)	PROG3P,311
CALL RANWT(DIV2,F1)	PROG3P,312
CALL RANRD(HUM1,F2)	PROG3P,313
CALL RANWT(HUM2,F2)	PROG3P,314
CALL RANRD(PSIM1,F2)	PROG3P,315
CALL RANWT(PSIM2,F2)	PROG3P,316
CALL RANRD(PSI11,F2)	PROG3P,317
CALL RANWT(PSI12,F2)	PROG3P,318
CALL RANRD(PSI21,F2)	PROG3P,319
CALL RANWT(PSI22,F2)	PROG3P,320
IF(IVAR.EQ.0) GO TO 70	PROG3P,321
IF(KIND.NE.0) GO TO 70	PROG3P,322
C INPUT OF BOUNDARY FIELD EACH SIX HOUR	PROG3P,323
C THE SUBROUTINE ZINPUT HAS TO BE WRITTEN ***	PROG3P,324
CALL RANRD(ZM2,F2)	PROG3P,325
CALL RANWT(ZM1,F2)	PROG3P,326
CALL RANRD(Z12,F2)	PROG3P,327
CALL RANWT(Z11,F2)	PROG3P,328
CALL RANRD(Z22,F2)	PROG3P,329
CALL RANWT(Z21,F2)	PROG3P,330
C CALL ZINPUT() ZPS,ZPM,ZP1 TO FIELDS F1,F2 AND F3,	PROG3P,331
DO 65 I=1,MN	PROG3P,332
Z=F2(I)	PROG3P,333
F2(I)=.5*(Z-F1(I))	PROG3P,334
F3(I)=.5*(F3(I)-Z)	PROG3P,335
F1(I)=Z	PROG3P,336
65 CONTINUE	PROG3P,337
CALL RANWT(ZM2,F1)	PROG3P,338
CALL RANWT(Z12,F2)	PROG3P,339
CALL RANWT(Z22,F3)	PROG3P,340
C GENERAL TIMESTEP THREE HOURS AHEAD	PROG3P,341
70 CALL STEP3P(F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,M,Y,MARK,IT,IVAR,M,N)	APR14,26
NCOUNT =NTIME	PROG3P,343
NTIME=NTIME+1	JUN12,11
I1=NTSTEP+1	PROG3P,345
C PRINT NUMBER OF ITERATIONS	PROG3P,346
WRITE(6,200) NCOUNT,NTIME	PROG3P,347
DO 71 I=1,I1	PROG3P,348
WRITE(6,201) IT(1,I),IT(2,I)	PROG3P,349
71 CONTINUE	PROG3P,350
C SMOOTHING AND ELLIPTICITY TEST	PROG3P,351
I1=0	PROG3P,352
I2=0	PROG3P,353
DO 80 I=1,10	PROG3P,354
IF(ABS(FLOAT(NSMUTT(I))-NTIME).LT.0.1) I1=1	JUN12,12
IF(ABS(FLOAT(MSMUTT(I))-NTIME).LT.0.1) I1=1	JUN12,13
IF(ABS(FLOAT(MELLIPT(I))-NTIME).LT.0.1) I2=1	JUN12,14
IF(ABS(FLOAT(NELLIPT(I))-NTIME).LT.0.1) I2=1	JUN12,15
80 CONTINUE	PROG3P,357
IF(I1.EQ.0.AND.I2.EQ.0) GO TO 95	PROG3P,358
CALL RANPD(PSIM1,F1)	PROG3P,359
CALL RANRD(PSI11,F2)	PROG3P,360
CALL RANRD(PSI21,F3)	PROG3P,361
DO 85 I=1,MN	PROG3P,362
Z=F1(I)	PROG3P,363

PROG3P (Continued)

F1(I)=Z+2*F2(I)	PROG3P,364
F3(I) = Z+2.*F3(I)	PROG3P,365
F2(I)=Z	PROG3P,366
85 CONTINUE	PROG3P,367
IF(I1.EQ.0) GO TO 86	PROG3P,368
CALL ASMUT(F1,F4,M,N,.5)	PROG3P,369
CALL ASMUT(F1,F4,M,N,.5)	PROG3P,370
CALL ASMUT(F2,F4,M,N,.5)	PROG3P,371
CALL ASMUT(F2,F4,M,N,.5)	PROG3P,372
CALL ASMUT(F3,F4,M,N,.5)	PROG3P,373
CALL ASMUT(F3,F4,M,N,.5)	PROG3P,374
86 IF(I2.EQ.0) GO TO 90	PROG3P,375
IF(KIND.NE.0) GO TO 90	APR14,27
CALL ELLIPT(F1,MY,M,N)	APR14,28
CALL ELLIPT(F2,MY,M,N)	APR14,29
CALL ELLIPT(F3,MY,M,N)	PROG3P,380
90 DO 93 I=1,MN	PROG3P,381
Z=F2(I)	PROG3P,382
F2(I)=.5*(Z-F1(I))	PROG3P,383
F3(I)=.5*(F3(I)-Z)	PROG3P,384
F1(I)=Z	PROG3P,385
93 CONTINUE	PROG3P,386
CALL RANWT(PSIM1,F1)	PROG3P,387
CALL RANWT(PSI11,F2)	PROG3P,388
CALL RANWT(PSI21,F3)	PROG3P,389
C MAPPRINTING	PROG3P,390
95 DO 96 I=1,10	PROG3P,391
IF(MAPHOUR(I).EQ.NTIME) GO TO 97	PROG3P,392
96 CONTINUE	PROG3P,393
GO TO 100	PROG3P,394
97 IF(KIND.NE.0) GO TO 98	APR14,30
CALL MAP3P(1,1,1,0,1,1,1,0,0,0,F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,	PROG3P,396
X MARK,M,N,IDAY,ITIME,NTIME,PNIVS,PNIVM,PNIV1,ZB,FB,SB)	PROG3P,397
GO TO 100	APR14,31
98 CALL MAP3P(0,0,0,0,1,0,0,0,1,1,0,F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,	PROG3P,399
X MARK,M,N,IDAY,ITIME,NTIME,PNIVS,PNIVM,PNIV1,ZB,ZB,ZB)	PROG3P,400
C ZERO TO ACCUMULATED PRECIPITATION	PROG3P,401
100 DO 105 I=1,10	PROG3P,402
IF(LETACC(I).EQ.NTIME) GO TO 106	PROG3P,403
105 CONTINUE	PROG3P,404
GO TO 115	PROG3P,405
106 DO 110 I=1,MN	PROG3P,406
F1(I)=0.0	PROG3P,407
110 CONTINUE	PROG3P,408
CALL RANWT(PREC,F1)	PROG3P,409
200 FORMAT(29H,NUMBER OF ITERATIONS BETWEEN,13.4H AND,13.6H HOURS)	PROG3P,410
201 FORMAT(4X,2(I4))	PROG3P,411
115 IF(MEND.LE.NTIME) STOP	PROG3P,412
GO TO 60	PROG3P,413
END	

SUBROUTINE STEP3P

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SUBROUTINE STEP3P(F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,MY,MARK,IT,
X          IVAR,M,N)
C SIX HOURS FORECAST BY 3P-MODEL
DIMENSION F1(1),F2(1),F3(1),F4(1),F5(1),F6(1),F7(1),F8(1),F9(1),
X          F10(1),MY(1),F(1),MARK(1),IT(2,1)
COMMON/FPAR/IC(8),JC(8),XPOL,YPOL,R,RE,DS,JMIN(100),JMAX(100),
X          MX,NX,KIND
COMMON/ECS/ PSIM1,PSI11,PSI21,PSIM2,PSI12,PSI22,HUM1,HUM2,DIV1,
2DIV2,WS,HEAT,J789,J12,J56,J3,PS,TS,PREC,STRM,STR1,STR2,ZM1,Z11,Z21
3,ZM2,Z12,Z22,H13,H23,HM3,HM2,H12,H22,H11,H21,J4,VM,V1,V2
COMMON/COEFF/A1,A2,A3,B1,B2,B3,C1,C2,C3,C4,C5,C6,C7,C8,D,DELP,EM,
X          E1,E2,H1,H2,H3,H4,H5,H6,PMEAN,S1,S2,T1,T2,T3,T4,T5,
X          P0,PM,P1
COMMON/COEFF2/T6,T7,T8,T9,T10,T11,T12,T13,T14,
X          K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15,
X          K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28,
X          K29,K30,K31,K32,K33,K34,K35,K36,K37,K38
COMMON/RUNPAR/DELT,NTSTEP,ALFASYS,ALFAM,ALFAZ,ALFAPSI,RESSYS,RESM,
X          RESZ,RESPSI,Q,FOCEAN,FCONT,WGT1,WGT2,WGT3,ADIFF
COMMON F
REAL MY,KEFF
REAL          K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15,
X          K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28,
X          K29,K30,K31,K32,K33,K34,K35,K36,K37,K38
DATA RGAS,EE,HL,CP,T0,E0,DEL1,DEL2,TOL/287,...622,2.5E6,1004,,
X          273,...611,0.0,0.0,0.0,0.0/
C DEL1,DEL2, AND TOL HAVE TO BE DEFINED *****
DEL1=0.0001
DEL2=0.001
TOL=0.0
C CONSTANTS FOR COMPUTATION OF LATENT HEAT
CC1 = 1./T0
CC2 = EE*HL/RGAS
CC3 = EE*HL/CP
CC4 = CC2*CC3
CC5 = DEL1+DEL2
C
IF(KIND.EQ.0) GO TO 5
WGT1=1.0
WGT2=.67
WGT3=.33
5 MN=M*N
ND=NTSTEP+1
M1=M-1
DO 170 KT=1,ND
EPS = 2.
IF(KT.LT.3) EPS = .5*KT
C
C*****JACOBIAN COMPUTATIONS*****
C ALL JACOBIANS ARE COMPUTED AND STORED
CALL RANRD(PSIM1,F1)
CALL RANRD(PSI11,F2)
CALL RANRD(PSI21,F3)
C
CALL ABSVOR(F1,F4,MY,MARK,M,N)
CALL RELVOR(F2,F5,MY,MARK,M,N)
CALL RELVOR(F3,F6,MY,MARK,M,N)
KEFF=1.0
CALL RANRD(PS,MY)
CALL DLOG(F5,MY,M,N,KEFF)
C

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STEP3P (Continued)

DO 10 I=1,MN	STEP3P.71
F7(I) = C3*F4(I)+C2*F5(I)+C4*F6(I) + C7*F(I)	STEP3P.72
F8(I) = C2*F4(I)+C5*F5(I)	STEP3P.73
F9(I) = C4*F4(I) +C6*F6(I)+2*C7*F(I)	STEP3P.74
10 CONTINUE	STEP3P.75
C CALL JACOB(F1,F7,F10,M,N)	STEP3P.76
CALL JACOB(F2,F8,F7,M,N)	STEP3P.77
CALL JACOB(F3,F9,F8,M,N)	STEP3P.78
CALL OROG(F7,MY,M,N,KEFF)	STEP3P.79
C	STEP3P.80
DO 20 I=1,MN	STEP3P.81
F10(I) = F10(I)+F8(I)+F7(I)	STEP3P.82
F9(I) = F4(I)+2*F5(I)	STEP3P.83
20 CONTINUE	STEP3P.84
C CALL RANWT(J789,F10)	STEP3P.85
C	STEP3P.86
CALL JACOB(F2,F9,F7,M,N)	STEP3P.87
CALL OROG(F7,MY,M,N,KEFF)	STEP3P.88
CALL JACOB(F1,F5,F8,M,N)	STEP3P.89
C	STEP3P.90
DO 30 I=1,MN	STEP3P.91
F10(I) = F7(I)+F8(I)	STEP3P.92
F9(I) = F4(I)+2*F6(I)	STEP3P.93
30 CONTINUE	STEP3P.94
C	STEP3P.95
CALL RANWT(J12,F10)	STEP3P.96
C	STEP3P.97
CALL JACOB(F3,F9,F7,M,N)	STEP3P.98
CALL JACOB(F1,F6,F8,M,N)	STEP3P.99
CALL JACOB(F1,F3,F9,M,N)	STEP3P.100
CALL JACOB(F1,F2,F10,M,N)	STEP3P.101
CALL OROG(F10,MY,M,N,KEFF)	STEP3P.102
C	STEP3P.103
CALL RANWT(J3,F10)	STEP3P.104
C	STEP3P.105
DO 40 I=1,MN	STEP3P.106
40 F10(I)=F7(I)+F8(I)	STEP3P.107
CALL RANWT(J56,F10)	STEP3P.108
CALL RANRD(PS,F10)	STEP3P.109
PPM=2./(P0-PM)	STEP3P.110
DO 44 I=1,MN	STEP3P.111
F7(I)=F1(I)-F2(I)*PPM*(F10(I)-PM)	STEP3P.112
F8(I)=PPM*F5(I)*(F10(I)-PM)-F4(I)+F(I)	STEP3P.113
44 F4(I)=F10(I)	STEP3P.114
C	STEP3P.115
CALL JACOB(F7,F4,F5,M,N)	STEP3P.116
CALL MYFF(MY,M,N)	STEP3P.117
C	STEP3P.118
C*****LOWER BOUNDARY CONDITION*****	APR14.51
C INFLUENCE FROM TOPOGRAPHY AND FRICTIO OVERR LAND OR OCEAN SURFACE	STEP3P.120
C OCEAN SURFACE IS ASSUMED WHERE STANDARD PRESSURE PS IS 101.35 CB OR MOST	STEP3P.121
DO 50 I=1,MN	STEP3P.122
IF(MARK(I)) 49,50,50	STEP3P.123
49 CF=FOCEAN	STEP3P.124
PP=F4(I)	STEP3P.125
IF(PP,LT,101.35) CF=FCONT	STEP3P.126
F4(I) = .25*MY(I)*F5(I)+CF*F8(I)	STEP3P.127
F10(I)=PP	STEP3P.128
50 CONTINUE	STEP3P.129
	STEP3P.130
	STEP3P.131

STEP3P (Continued)

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C
CALL RANWT(WS,F4)
CALL RANRD(TS,F5)
C*****SENSIBLE HEAT*****
C HEATING FROM OCEAN SURFACE WHEN THE AIR IS COLDER, PQ IS THE TEMP DIFF
C OCEAN SURFACE IS ASSUMED WHERE STANDARD PRESSURE PS IS 101.35 CB OR MOST
DO 59 I=2,M1
J1=JMIN(I)+1
J2=JMAX(I)-1
K=(J1-1)*M+1
DO 59 J=J1,J2
PP=F10(K)
PQ=F5(K)-H2*F2(K)
IF(PP.LT.101.324) GO TO 57
IF(PQ.LE.0.0) GO TO 57
PR = SQRT(.25*MY(K)*((F7(K+1)-F7(K-1))*(F7(K+1)-F7(K-1))+
X (F7(K+M)-F7(K-M))*(F7(K+M)-F7(K-M))))
F5(K) = .5E-2*H1*(.1*PR-1.)*PQ
GO TO 58
57 F5(K)=0.0
58 K=K+M
59 CONTINUE
C*****HUMIDITY FORECAST*****
C EM,E1,E2 ARE COEFF FOR COMP OF MEAN STREAMFUNCTION
C SS1,SS2,SS3 ARE COEFF FOR COMP OF MEAN DIVERGENCE, D IS ZERO OVER LAND
C THE HUMIDITY IS GIVEN IN TON PER SQUAREMETER AND CENTIBAR
CALL RANRD(DIV1,F6)
CALL RANRD(DIV2,F7)
CALL RANRD(HUM1,F8)
C
SS1 = E1+EM*C2
SS2 = E2-EM*C1
SS3 = -EM*C3
DO 61 I=1,MN
PQ = D
PP = F10(I)
IF(PP.LT.101.35) PQ=0.0
F7(I) = F8(I)*(SS1*F6(I)+SS2*F7(I)+F4(I)*(PQ+SS3))
F6(I) = EM*F1(I)+E1*F2(I)+E2*F3(I)
61 CONTINUE
CALL JACOB(F6,F8,F10,M,N)
DO 62 I=2,M1
J1=JMIN(I)+1
J2=JMAX(I)-1
K=(J1-1)*M+1
DO 62 J=J1,J2
F6(K) = F8(K+M)+F8(K-M)+F8(K+1)+F8(K-1)-4*F8(K)
K=K+M
62 CONTINUE
C
DEPS = EPS*DELT
DO 63 I=1,MN
IF(MARK(I).GE.0) GO TO 63
F8(I)=-DEPS*(.25*MY(I)*F10(I)+F7(I)-ADIFF*MY(I)*F6(I))
63 CONTINUE
C
CALL RANRD(HUM2,F6)
DO 65 I=1,MN
IF(MARK(I)) 64,65,65
64 F6(I) = F6(I)+F8(I)
65 CONTINUE
C*****PRECIPITATION*****

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STEP3P (Continued)

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C THE RAIN FOR ONE TIMESTEP IS ACCUMULATED. THE RAIN IS GIVEN IN MM OR KSTEP3P.194
C PER SQUAREMETER - STEP3P.195
DO 70 I=1,MN STEP3P.196
IF(MARK(I)) 66,70,70 STEP3P.197
66 TEMP = H3*F2(I)+H4*F3(I) STEP3P.198
CALL SATUR(TEMP,QSAT) STEP3P.199
PQ = F6(I)-.8*QSAT STEP3P.200
IF(PQ) 68,68,67 STEP3P.201
67 F6(I) = .8*QSAT STEP3P.202
F7(I) = .5E3*DELP*PQ STEP3P.203
IF(KT,EQ.1) F7(I)=0,0 STEP3P.204
IF(KT,EQ.2) F7(I)=2*F7(I) STEP3P.205
GO TO 70 STEP3P.206
68 F7(I) = 0 STEP3P.207
PQ = F6(I)-.2*QSAT STEP3P.208
IF(PQ) 69,70,70 STEP3P.209
69 F6(I) = .2*QSAT STEP3P.210
70 CONTINUE STEP3P.211
CALL RANRD(PREC,F8) STEP3P.212
DO 71 I=1,IN STEP3P.213
F8(I) = F8(I) + F7(I) STEP3P.214
71 CONTINUE STEP3P.215
CALL RANNT(PREC,F8) STEP3P.216
CALL RANRD(HUM1,F8) STEP3P.217
CALL RANNT(HUM1,F6) STEP3P.218
CALL RANNT(HUM2,F8) STEP3P.219
C*****LATENT HEAT*****STEP3P.220
CALL RANRD(DIV1,F6) STEP3P.221
CALL RANRD(DIV2,F8) STEP3P.222
DO 180 I=1,MN STEP3P.223
IF(MARK(I)) 179,180,180 STEP3P.224
179 F6(I) = T1*F4(I)+T2*F6(I)+T3*F8(I) STEP3P.225
180 CONTINUE STEP3P.226
C STEP3P.227
DO 190 I=1,MN STEP3P.228
IF(MARK(I)) 187,190,190 STEP3P.229
187 RAIN = F7(I)/EPS/DELT STEP3P.230
IF(RAIN,LT,TOL) GO TO 188 STEP3P.231
VERT = F6(I) STEP3P.232
IF(VERT,GT,-DEL1) GO TO 188 STEP3P.233
OSTAR = VERT STEP3P.234
IF(VERT,GE,-CC5,AND,VERT,LE,-DEL1) OSTAR = -ABS(VERT*VERT/CC5) STEP3P.235
TEMP = H6*F2(I) STEP3P.236
X = CC2*(CC1-1./TEMP) STEP3P.237
E = E0*EXP(X) STEP3P.238
FSTAR = EE*TEMP*E*(CC3-TEMP)/PMEAN/(PMEAN*TEMP*TEMP + CC4*E) STEP3P.239
HLAT = -H1*HL*OSTAR*FSTAR STEP3P.240
GO TO 189 STEP3P.241
188 HLAT = 0,0 STEP3P.242
189 F6(I) = F6(I) + HLAT STEP3P.243
190 CONTINUE STEP3P.244
CALL RANNT(HEAT,F5) STEP3P.245
C STEP3P.246
I1=0 STEP3P.247
I2=0 STEP3P.248
I3=0 STEP3P.249
I4=0 STEP3P.250
CALL STEPEXT(F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,MY,MARK,M,N, APR14.56
X I1,I2,I3,I4) STEP3P.253
C*****FORCING FUNCTIONS*****STEP3P.254

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STEP3P (Continued)

CALL RANRD(J789,F10)	STEP3P.255
DEPS = .25*DELT*EPS	STEP3P.256
DO 73 I=1,MN	STEP3P.257
IF(MARK(I)) 72,73,73	STEP3P.258
72 PQ = 4*F(I)/MY(I)	STEP3P.259
F8(I) =-DEPS*(F10(I) - C8*F4(I)*PQ)	STEP3P.260
F6(I) = PQ*(A3*F4(I)+A1*F5(I))	STEP3P.261
F7(I)=PQ*(B3*F4(I)+B1*F5(I))	STEP3P.262
73 CONTINUE	STEP3P.263
C	STEP3P.264
CALL RANRD(J3,F10)	STEP3P.265
CALL RANRD(J12,F4)	STEP3P.266
CALL RANRD(J56,F5)	STEP3P.267
C	STEP3P.268
DO 80 I=1,MN	STEP3P.269
IF(MARK(I)) 79,80,80	STEP3P.270
79 PQ = F(I)*F(I)	STEP3P.271
F6(I) =-DEPS*(F4(I)+PQ*(A2*F9(I)-A1*F10(I))+F6(I))	STEP3P.272
F7(I) =-DEPS*(F5(I)+PQ*(B1*F10(I)-B2*F9(I))+F7(I))	STEP3P.273
80 CONTINUE	STEP3P.274
C*****SOLUTION OF FORECAST FQ,*****	STEP3P.275
CALL RANRD(HM3,F4)	STEP3P.276
CALL RANRD(H13,F5)	STEP3P.277
CALL RANRD(H23,F10)	STEP3P.278
DEPS1=DELT*EPS	STEP3P.279
DO 81 I=1,MN	STEP3P.280
F8(I)=F8(I)+DEPS1/MY(I)*F4(I)	STEP3P.281
F6(I)=F6(I)+DEPS1/MY(I)*F5(I)	STEP3P.282
81 F7(I)=F7(I)+DEPS1/MY(I)*F10(I)	STEP3P.283
CALL RANRD(PS112,F5)	STEP3P.284
CALL RANRD(PS122,F10)	STEP3P.285
C	STEP3P.286
DO 90 I=1,MN	STEP3P.287
89 F2(I)=2*(F2(I)-F5(I))	STEP3P.288
F3(I)=2*(F3(I)-F10(I))	STEP3P.289
90 CONTINUE	STEP3P.290
C	STEP3P.291
LABEL = 10H HELMSYS	APR14.57
BTIME = SECOND(DUMMY)	APR14.58
WRITE(6,8000) BTIME	APR14.59
8000 FORMAT(1H , *BTIME= *, F10.4)	APR14.60
CALL HELISYS(F2,F3,F6,F7,MY,F4,A1,A2,B1,B2,ALFASYS,RESSYS,	APR14.61
X ITSYS,M,N)	STEP3P.293
DTIME = BTIME - SECOND(DUMMY)	APR14.62
WRITE(6,8005) LABEL, DTIME	APR14.63
8005 FORMAT(1H , *TIME TO EXECUTE *, A10, F10.4)	APR14.64
IT(1,KT) = ITSYS	STEP3P.294
CALL ASMUT(F2,F4,M,N,,5)	JUN12.18
CALL ASMUT(F2,F4,M,N,-,5)	JUN12.19
CALL ASMUT(F3,F4,M,N,,5)	JUN12.20
CALL ASMUT(F3,F4,M,N,-,5)	JUN12.21
C	STEP3P.295
DO 100 I=1,MN	STEP3P.296
IF(MARK(I)) 99,100,100	STEP3P.297
99 TFILT=0.4	JUN12.22
IF(MARK(I).EQ.-10) TFILT=0.7	JUN12.23
IF(MARK(I).EQ.-1) TFILT=1.	JUN12.24
F5(I)=F5(I) + TFILT*F2(I)	JUN12.25
F10(I)=F10(I) + TFILT*F3(I)	JUN12.26
F4(I) = 0/MY(I)	STEP3P.300
100 CONTINUE	STEP3P.301
C	STEP3P.302

STEP3P (Continued)

C	CALL RANRD(PSIM2,F6)	STEP3P.303
	DO 110 I=1,MN	STEP3P.304
109	F1(I) = 2*(F1(I)-F6(I)) -C2*F2(I) +C1*F3(I)	STEP3P.305
110	CONTINUE	STEP3P.306
C	CALL HELM(F1,F8,F4,ALFAM,RESM,ITM,M,N)	STEP3P.307
	IT(2,KT) = ITM	STEP3P.308
C	CALL ASMUT(F1,F4,M,N,.5)	STEP3P.309
	CALL ASHUT(F1,F4,M,N,-.5)	STEP3P.310
	DO 120 I=1,MN	STEP3P.311
	IF(MARK(I))119,120,120	JUN12,27
119	TFILT=.4	JUN12,28
	IF(MARK(I).EQ.-10) TFILT=0.7	STEP3P.312
	IF(MARK(I).EQ.-1) TFILT=1.	JUN12,29
	F6(I)=F6(I) + TFILT*F1(I) + C2*F2(I)	JUN12,30
	S=C1*F3(I)	JUN12,31
120	CONTINUE	JUN12,32
C	*****MIXING WITH BOUNDARY FIELDS*****	JUN12,33
	CALL RANRD(STRM,F4)	JUN12,34
	IF(IVAR.EQ.0) GO TO 156	STEP3P.314
	IF(KIND.NE.0) GO TO 156	STEP3P.315
	CALL RANPD(ZM1,F7)	STEP3P.316
	CALL RANRD(ZM2,F8)	STEP3P.317
C	WF = (KT-1)/ND	STEP3P.318
	FW = 1.-WF	STEP3P.319
	G = 9.806	STEP3P.320
	DO 130 I=1,MN	STEP3P.321
	IF(MARK(I)) 129,130,130	STEP3P.322
129	F7(I) = G*(FW*F7(I)+WF*F8(I))/F(I)	STEP3P.323
130	CONTINUE	STEP3P.324
	CALL MIXF(F6,F7,MARK,WGT1,WGT2,WGT3,M,N)	STEP3P.325
	CALL RANRD(STR1,F4)	STEP3P.326
C	CALL RANRD(Z11,F7)	JUN12,35
	CALL RANRD(Z12,F8)	STEP3P.328
	DO 140 I=1,MN	STEP3P.329
	IF(MARK(I)) 139,140,140	STEP3P.330
139	F7(I) = G*(FW*F7(I)+WF*F8(I))/F(I)	STEP3P.331
140	CONTINUE	STEP3P.332
	CALL MIXF(F5,F7,MARK,WGT1,WGT2,WGT3,M,N)	STEP3P.333
	CALL RANRD(STR2,F4)	STEP3P.334
C	CALL RANRD(Z21,F7)	STEP3P.335
	CALL RANRD(Z22,F8)	JUN12,36
	DO 150 I=1,MN	STEP3P.337
	IF(MARK(I)) 149,150,150	STEP3P.338
149	F7(I) = G*(FW*F7(I)+WF*F8(I))/F(I)	STEP3P.339
150	CONTINUE	STEP3P.340
	CALL MIXF(F10,F7,MARK,WGT1,WGT2,WGT3,M,N)	STEP3P.341
	GO TO 156	STEP3P.342
151	CALL MIXF(F6,F4,MARK,WGT1,WGT2,WGT3,M,N)	STEP3P.343
	CALL RANRD(STRM,F4)	STEP3P.344
	CALL MIXF(F5,F4,MARK,WGT1,WGT2,WGT3,M,N)	JUN12,37
	CALL RANRD(STR1,F4)	STEP3P.346
	CALL MIXF(F10,F4,MARK,WGT1,WGT2,WGT3,M,N)	STEP3P.347
	CALL RANRD(STR2,F4)	STEP3P.348
C	*****STORE NEW TIMESTEP*****	STEP3P.349

STEP3P (Continued)

156	IF(KT.LT.3) GO TO 157	STEP3P,356
	CALL RANRD(PSIM1,F1)	STEP3P,357
	CALL RANRD(PSI11,F4)	STEP3P,358
	CALL RANRD(PSI21,F7)	STEP3P,359
157	CALL RANWT(PSIM1,F6)	STEP3P,360
	CALL RANWT(PSI11,F5)	STEP3P,361
	CALL RANWT(PSI21,F10)	STEP3P,362
	IF(KT.LT.3) GO TO 158	STEP3P,363
	CALL RANWT(PSIM2,F1)	STEP3P,364
	CALL RANWT(PSI12,F4)	STEP3P,365
	CALL RANWT(PSI22,F7)	STEP3P,366
C	*****COMPUTATION OF DIVERGENCE*****	STEP3P,367
158	CALL RANRD(J3,F10)	STEP3P,368
	CALL RANRD(HEAT,F5)	STEP3P,369
	CALL RANRD(WS,F8)	STEP3P,370
C		STEP3P,371
	PQ = 1./EPS/DELT	STEP3P,372
	DO 160 I=1,MN	STEP3P,373
	IF(MARK(I)) 159,160,160	STEP3P,374
159	TERM1 = F(I)*(PQ *F2(I)+.25*MY(I)*F10(I))-F5(I)-S1 *F8(I)	STEP3P,375
	TERM2 = F(I)*(PQ *F3(I)+.25*MY(I)*F9 (I)) -S2 *F8(I)	STEP3P,376
	F2(I) =- A1*TERM1 + A2*TERM2	STEP3P,377
	F3(I) =- B1*TERM1 - B2*TERM2	STEP3P,378
160	CONTINUE	STEP3P,379
C		STEP3P,380
	CALL RANWT(DIV1,F2)	STEP3P,381
	CALL RANWT(DIV2,F3)	STEP3P,382
	CALL RANRD(PSIM1,F1)	STEP3P,383
	CALL RANRD(PSI11,F5)	STEP3P,384
	CALL RANRD(PSI21,F10)	STEP3P,385
C		STEP3P,386
	PRINT 7512,KT	STEP3P,387
7512	FORMAT(1X,3HKT=,I3)	STEP3P,388
170	CONTINUE	STEP3P,389
	RETURN	STEP3P,390
	END	STEP3P,391

SUBROUTINE STEPEXT

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SUBROUTINE STEPEXT(F1,F2,F3,F4,F5,F6,F7,F8,F9,F10,MY,MARK,M,N, APR14,67
1 11,I2,I3,I4) STEPEXT,3
C THIS SUBROUTINE COMPUTES THE CONTRIBUTION FROM THE HIGHER ORDER TERMS STEPEXT,4
C IN THE VORTICITY EQUATION, STEPEXT,5
C I1 INDICATES THE VORTICITY ADVECTION BY THE DIVERGENT VIND STEPEXT,6
C I2 INDICATES THE RELATIVE VORTICITY*DIVERGENCE STEPEXT,7
C I3 INDICATES THE VERTICAL ADVECTION OF VORTICITY STEPEXT,8
C I4 INDICATES THE TWISTINGTERM STEPEXT,9
C I1=0 NO CONTRIBUTION I1 DIFFERENT FROM 0 CONTRIBUTION FROM I1 STEPEXT,10
C I2=0 NO CONTRIBUTION I2 DIFFERENT FROM 0 CONTRIBUTION FROM I2 STEPEXT,11
C I3=0 NO CONTRIBUTION I3 DIFFERENT FROM 0 CONTRIBUTION FROM I3 STEPEXT,12
C I4=0 NO CONTRIBUTION I4 DIFFERENT FROM 0 CONTRIBUTION FROM I4 STEPEXT,13
C PSIM IS IN F1,PSI1 IS IN F2,PSI2 IS IN F3,WS IS IN F4 STEPEXT,14
C STEPEXT NEEDS 10 FIELDS IN THE FAST CORE MEMORY F,MY AND MARK MUST STEPEXT,15
C ALSO BE IN FAST CORE MEMORY STEPEXT,16
DIMENSION F1(1),F2(1),F3(1),F4(1),F5(1),F6(1),F7(1),F8(1),F9(1), STEPEXT,17
1 F10(1),F(1),MARK(1),MY(1) STEPEXT,18
COMMON F APR14,68
COMMON/COEFF/A1,A2,A3,B1,B2,B3,C1,C2,C3,C4,C5,C6,C7,C8,D,DELP,EM, STEPEXT,19
1 E1,E2,H1,H2,H3,H4,H5,H6,PMEAN,S1,S2,T1,T2,T3,T4,T5, STEPEXT,20
2 P0,PM,P1 STEPEXT,21
COMMON/COEFF2/T6,T7,T8,T9,T10,T11,T12,T13,T14, STEPEXT,22
X K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15, STEPEXT,23
X K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28, STEPEXT,24
X K29,K30,K31,K32,K33,K34,K35,K36,K37,K38 STEPEXT,25
COMMON/ECS/ PSIM1,PSI11,PSI21,PSIM2,PSI12,PSI22,HUM1,HUM2,DIV1, APR14,69
2 DIV2,WS,HEAT,J789,J12,J56,J3,PS,TS,PREG,STRM,STR1,STR2,ZH1,Z11,Z21 APR14,70
3,ZM2,Z12,Z22,H13,H23,H43,H42,H12,H22,H11,H21,J4,VM,V1,V2 APR14,71
REAL MY STEPEXT,34
REAL K1,K2,K3,K4,K5,K6,K7,K8,K9,K10,K11,K12,K13,K14,K15, STEPEXT,35
X K16,K17,K18,K19,K20,K21,K22,K23,K24,K25,K26,K27,K28, STEPEXT,36
X K29,K30,K31,K32,K33,K34,K35,K36,K37,K38 STEPEXT,37
KIND=0 APR14,72
MN=1+M STEPEXT,38
RESIDUE =.5E4 STEPEXT,39
ALFA=1.4 STEPEXT,40
C JACOBIAN J4 TO SECONDARY STORAGE,DIVERGENCIES TO FAST MEMORY STEPEXT,41
CALL RANHT(J4,F9) STEPEXT,42
CALL RANHD(DIV1,F5) STEPEXT,43
CALL RANHD(DIV2,F6) STEPEXT,44
DO 9 I=1,MN STEPEXT,45
IF(MARK(I)) 9,7,7 STEPEXT,46
7 F4(I)=0.0 STEPEXT,47
F5(I)=0.0 STEPEXT,48
F6(I)=0.0 STEPEXT,49
9 CONTINUE STEPEXT,50
IF(KIND.EQ.0) GO TO 8 STEPEXT,51
CALL BMOVE(F4,M,N) STEPEXT,52
CALL BMOVE(F5,M,N) STEPEXT,53
CALL BMOVE(F6,M,N) STEPEXT,54
8 CONTINUE STEPEXT,55
IF(I4.EQ.0.AND.I3.EQ.0.AND.I2.EQ.0.AND.I1.EQ.0) GO TO 170 STEPEXT,56
IF(I4.EQ.0) GO TO 44 STEPEXT,57
C COMPUTE THE TWISTINGTERM,I4 STEPEXT,58
DO 10 I=1,MN STEPEXT,59
F7(I)=K31*F5(I)+K32*F6(I)+K33*F4(I) STEPEXT,60
10 F8(I)=K36*F5(I)+K37*F6(I)+K38*F4(I) STEPEXT,61
CALL GRADPR(F2,F7,F9,MARK,M,N) STEPEXT,62
CALL GRADPR(F3,F8,F10,MARK,M,N) STEPEXT,63
DO 11 I=1,MN STEPEXT,64
F9(I)=0.5*MY(I)*F9(I) STEPEXT,65

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STEPEXT (Continued)

11 F10(I)=0.5*MY(I)*F10(I)	STEPEXT.66
CALL RANWT(H13,F9)	STEPEXT.67
CALL RANWT(H23,F10)	STEPEXT.68
DO 20 I=1,MN	STEPEXT.69
F7(I)=K19*F5(I)+K20*F6(I)+K21*F4(I)	STEPEXT.70
20 F8(I)=K22*F5(I)+K23*F6(I)+K24*F4(I)	STEPEXT.71
CALL GRADPR(F2,F7,F9,MARK,M,N)	STEPEXT.72
CALL GRADPR(F3,F8,F10,MARK,M,N)	STEPEXT.73
DO 30 I=1,MN	STEPEXT.74
30 F7(I)=K25*F5(I)+K26*F6(I)+K27*F4(I)	STEPEXT.75
CALL GRADPR(F1,F7,F8,MARK,M,N)	STEPEXT.76
DO 40 I=1,MN	STEPEXT.77
40 F8(I)=(F8(I)+F9(I)+F10(I))*0.5*MY(I)	STEPEXT.78
CALL RANWT(HM3,F8)	STEPEXT.79
GO TO 45	STEPEXT.80
44 DO 46 I=1,MN	STEPEXT.81
46 F8(I)=0.0	STEPEXT.82
CALL RANWT(H13,F8)	STEPEXT.83
CALL RANWT(H23,F8)	STEPEXT.84
CALL RANWT(HM3,F8)	STEPEXT.85
45 IF(I2.EQ.0.AND.I3.EQ.0.AND.I1.EQ.0) GO TO 165	STEPEXT.86
50 CALL RELVOR(F1,F7,MY,MARK,M,N)	APR14.73
CALL RELVOR(F2,F8,MY,MARK,M,N)	APR14.74
CALL RELVOR(F3,F9,MY,MARK,M,N)	APR14.75
IF(I2.EQ.0.AND.I3.EQ.0) GO TO 94	STEPEXT.90
IF(I2.EQ.0) GO TO 65	STEPEXT.91
C COMPUTATION OF THE RELATIVE VORTICITY*DIVERGENCE,I2	STEPEXT.92
DO 60 I=1,MN	STEPEXT.93
F1(I)=F5(I)*(K1*F8(I)+K2*F9(I)+K3*F7(I))+F6(I)*(K4*F8(I)+K5*F9(I)+	STEPEXT.94
1 K6*F7(I))+F4(I)*(K7*F8(I)+K8*F9(I)+K9*F7(I))	STEPEXT.95
F2(I)=F5(I)*(K28*F8(I)+F7(I))+F6(I)*K29*F8(I)+F4(I)*K30*F8(I)	STEPEXT.96
60 F3(I)=F5(I)*K34*F9(I)+F6(I)*(K35*F9(I)+F7(I))+F4(I)*K30*F9(I)	STEPEXT.97
C COMPUTATION OF THE VERTICAL ADVECTION OF VORTICITY,I3	STEPEXT.98
IF(I3.EQ.0) GO TO 91	STEPEXT.99
GO TO 80	STEPEXT.100
65 DO 70 I=1,MN	STEPEXT.101
F1(I)=0.0	STEPEXT.102
F2(I)=0.0	STEPEXT.103
70 F3(I)=0.0	STEPEXT.104
80 DO 90 I=1,MN	STEPEXT.105
F1(I)=F1(I)+F5(I)*(K10*F8(I)+K11*F9(I)+K12*F7(I))+	STEPEXT.106
1 F6(I)*(K13*F8(I)+K14*F9(I)+K15*F7(I))+	STEPEXT.107
2 F4(I)*(K16*F8(I)+K17*F9(I)+K18*F7(I))	STEPEXT.108
F2(I)=F2(I)+F8(I)*(K31*F5(I)+K32*F6(I)+K33*F4(I))	STEPEXT.109
90 F3(I)=F3(I)+F9(I)*(K36*F5(I)+K37*F6(I)+K38*F4(I))	STEPEXT.110
GO TO 91	STEPEXT.111
94 DO 92 I=1,MN	STEPEXT.112
F1(I)=0.0	STEPEXT.113
F2(I)=0.0	STEPEXT.114
92 F3(I)=0.0	STEPEXT.115
91 CALL RANWT(HM2,F1)	STEPEXT.116
CALL RANWT(H12,F2)	STEPEXT.117
CALL RANWT(H22,F3)	STEPEXT.118
C COMPUTATION OF THE ADVECTION OF VORTICITY BY THE DIVERGENT WIND -I1	STEPEXT.119
C COMPUTE FORCINGFUNCTION FOR THE VELOCITYPOTENTIAL	STEPEXT.120
IF(I1.EQ.0) GO TO 166	STEPEXT.121
95 DO 100 I=1,MN	STEPEXT.122
F1(I)=(C2*F5(I)-C1*F6(I)-C8*F4(I))/MY(I)	STEPEXT.123
F2(I) = F5(I)/MY(I)	STEPEXT.124
100 F3(I) = F6(I)/MY(I)	STEPEXT.125
C SOLVE THE POISSONEQUATION BY RELAXATION IN ORDER TO GET VELOCITYPOT.	STEPEXT.126
CALL RANRD(VN,F4)	STEPEXT.127

STEPEXT (Continued)

CALL RANRD(V1,F5)	STEPEXT,128
CALL RANRD(V2,F6)	STEPEXT,129
CALL VELPOT(F4,F1,M,N,RESIDUE,ALFA)	STEPEXT,130
CALL VELPOT(F5,F2,M,N,RESIDUE,ALFA)	STEPEXT,131
CALL VELPOT(F6,F3,M,N,RESIDUE,ALFA)	STEPEXT,132
CALL RANWT(VM,F4)	STEPEXT,133
CALL RANWT(V2,F6)	STEPEXT,134
CALL RANWT(V1,F5)	STEPEXT,135
DO 110 I=1,MN	STEPEXT,136
F1(I)=F7(I)+F(I)-2.*F8(I)	STEPEXT,137
110 F2(I)=F7(I)+F(I)+2.*F9(I)	STEPEXT,138
CALL GRADPR(F4,F8,F3,MARK,M,N)	STEPEXT,139
CALL GRADPR(F5,F1,F10,MARK,M,N)	STEPEXT,140
DO 120 I=1,MN	STEPEXT,141
120 F10(I)=-0.5*MY(I)*(F3(I)+F10(I))	STEPEXT,142
CALL RANWT(H11,F10)	STEPEXT,143
CALL GRADPR(F4,F9,F3,MARK,M,N)	STEPEXT,144
CALL GRADPR(F6,F2,F10,MARK,M,N)	STEPEXT,145
DO 130 I=1,MN	STEPEXT,146
130 F10(I)=-0.5*MY(I)*(F3(I)+F10(I))	STEPEXT,147
CALL RANWT(H21,F10)	STEPEXT,148
DO 140 I=1,MN	STEPEXT,149
F1(I)=C3*(F7(I)+F(I))-C2*F8(I)+C4*F9(I)+C7*F(I)	STEPEXT,150
F2(I)=-C2*(F7(I)+F(I))+C5*F8(I)	STEPEXT,151
140 F3(I)=C4*(F7(I)+F(I))+C6*F9(I)+2.*C7*F(I)	STEPEXT,152
CALL GRADPR(F4,F1,F7,MARK,M,N)	STEPEXT,153
CALL GRADPR(F5,F2,F8,MARK,M,N)	STEPEXT,154
CALL GRADPR(F6,F3,F9,MARK,M,N)	STEPEXT,155
DO 150 I=1,MN	STEPEXT,156
150 F1(I)=-0.5*MY(I)*(F7(I)+F8(I)+F9(I))	APR14,76
CALL RANRD(HM2,F2)	STEPEXT,158
CALL RANRD(HM3,F3)	STEPEXT,159
CALL RANRD(H11,F4)	STEPEXT,160
CALL RANRD(H12,F5)	STEPEXT,161
CALL RANRD(H13,F6)	STEPEXT,162
CALL RANRD(H21,F7)	STEPEXT,163
CALL RANRD(H22,F8)	STEPEXT,164
CALL RANRD(H23,F9)	STEPEXT,165
DO 160 I=1,MN	STEPEXT,166
F1(I)=F1(I)+F2(I)+F3(I)	STEPEXT,167
F4(I)=F4(I)+F5(I)+F6(I)	STEPEXT,168
160 F7(I)=F7(I)+F8(I)+F9(I)	STEPEXT,169
GO TO 190	STEPEXT,170
163 CALL RANRD(HM3,F1)	STEPEXT,171
CALL RANRD(H13,F4)	STEPEXT,172
CALL RANRD(H23,F7)	STEPEXT,173
GO TO 190	STEPEXT,174
166 CALL RANRD(HM2,F2)	STEPEXT,175
CALL RANRD(HM3,F3)	STEPEXT,176
CALL RANRD(H12,F5)	STEPEXT,177
CALL RANRD(H13,F6)	STEPEXT,178
CALL RANRD(H22,F8)	STEPEXT,179
CALL RANRD(H23,F9)	STEPEXT,180
DO 167 I=1,MN	STEPEXT,181
F1(I)=F2(I)+F3(I)	STEPEXT,182
F4(I)=F5(I)+F6(I)	STEPEXT,183
167 F7(I)=F8(I)+F9(I)	STEPEXT,184
GO TO 190	STEPEXT,185
170 DO 180 I=1,MN	STEPEXT,186
F1(I)=0.0	STEPEXT,187
F4(I)=0.0	STEPEXT,188
180 F7(I)=0.0	STEPEXT,189

STEPEXT (Continued)

```
190 CALL RANWT(HM3,F1)
    CALL RANWT(H13,F4)
    CALL RANWT(H23,F7)
    CALL RANRD(PSIM1,F1)
    CALL RANRD(PSI11,F2)
    CALL RANRD(PSI21,F3)
    CALL RANRD(WS,F4)
    CALL RANRD(HEAT,F5)
    CALL RANRD(J4,F9)
200 RETURN
    END
```

```
STEPEXT,190
STEPEXT,191
STEPEXT,192
STEPEXT,193
STEPEXT,194
STEPEXT,195
STEPEXT,196
STEPEXT,197
STEPEXT,198
STEPEXT,199
STEPEXT,200
```


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